

Materia Oscura: Datos y Ideas

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en colaboración con

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Outline

- Evidence
- Not- evidence
- Within gravity
- Particles
- Light particles high energy signals
- Conclusion?

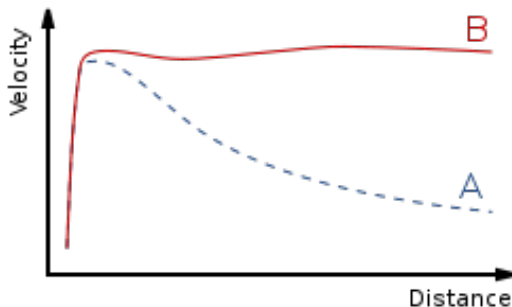


DM Evidence



DM Evidence

Galaxy rotation curves:



Stars outside are faster than visible mass inside would allow them to do



DM Evidence

Gravitational lensing, clusters of galaxies:

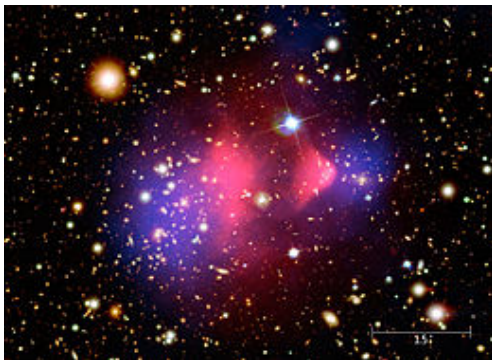


Something dark and transparent and massive
between galaxy cluster and us (DM).



DM Evidence

Gravitational lensing, collision of galaxies:

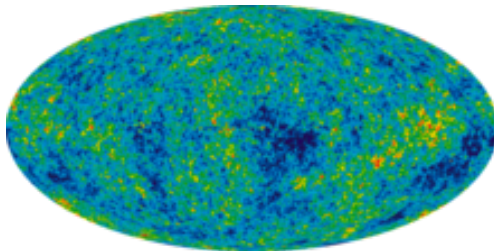


Something dark and transparent and massive in 2 galaxies.
But also tells the **DM** interacts weaker than interstellar gas.



DM Evidence

CMB fluctuations:

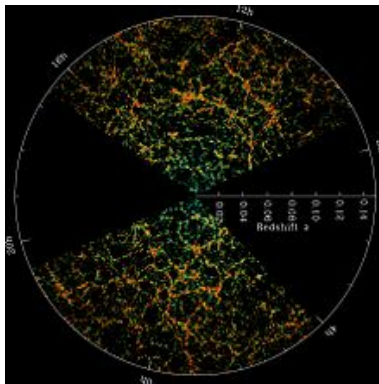


Escape from plasma \rightarrow neutral matter in recombination.
Only can model power spectrum of perturbations if one assumes certain amount of **dark energy and dark matter**.



DM Evidence

Barionic acoustic oscillations:

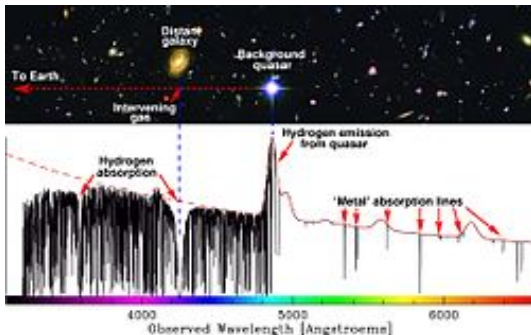


Remain from plasma \rightarrow neutral matter in recombination.
Only can model power spectrum of densities if one assumes certain
amount of **dark energy and dark matter**.



DM Evidence

Lyman alpha forrest:

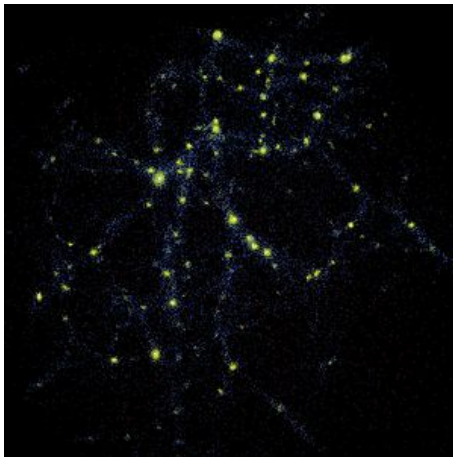


Model hydrogen absorption lines (without redshift 1216Å).
Only works with **dark energy and dark matter**



DM Evidence

Structure formation:



Numerical simulation of structure
needs also DM



DM Non-Evidence



DM Non-Evidence

Direct passive detection:



Many experiments, Xenon, Edelweiss, ...
Big detectors, low background, **no signal**



DM Non-Evidence

Direct passive detection:

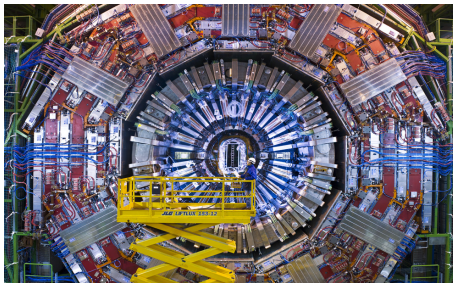


Many more experiments
Bigger detectors, lower background ... no signal?



DM Non-Evidence

Produce and detect:



Several experiments, ATLAS, CMS, ...

No signal



DM Non-Evidence

Indirect passive detection:



Several experiments, Fermi-LAT, CTA ...

No signal



What is the difference?

The non-detections involve SM interactions only

The yes-detections involve SM interactions **AND Gravity**

- Newton $U = G \frac{Mm}{r}$
- Newton $F = G \frac{Mm}{r^2}$ & geodesics $\frac{d^2 x^\mu}{dt^2} = \Gamma^\mu_{\alpha\beta} u^\alpha u^\beta$
- Universe evolution $H^2 = H_0^2(\Omega_M(1+z)^3 + \Omega_r(1+z)^2 + \Omega_\Lambda)$



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Within gravity



Within gravity

Maybe something wrong with Gravity
(remember evidence only from Newton and Hubble)

Ideas in this direction

- MOND
- Extra fields in EH action “Scalar-Tensor” ...
- Non-local operators $S = \int dx^4 R + R \frac{1}{\square} R + \dots$
- ...
- Our idea in this direction: “Scale dependence of gravitational couplings”*

*B.K. and Paola Riosco, arXiv:1501.00904;

D. Rodrigues, B.K., O. Piattella, I. Shapiro, AIP Conf.Proc. 1647 (2015) 57-61; D. Rodrigues, P. Letelier, I. Shapiro, arXiv:1102.2188;

JCAP 1004 (2010) 020.



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Within gravity

Scale dependence

Allow for “Scale dependence of gravitational couplings” modifies Einstein’s field equations.

De Sitter case:

eom $g_{\mu\nu}$:

$$G_{\mu\nu} = -g_{\mu\nu}\Lambda_k - \Delta t_{\mu\nu} \quad , \quad (1)$$

eom k :

$$R\nabla_\mu \left(\frac{1}{G_k} \right) - 2\nabla_\mu \left(\frac{\Lambda_k}{G_k} \right) = 0 \quad . \quad (2)$$

can be solved with “Schwarzschild Ansatz”



Within gravity

Scale dependence

generalized de Sitter solution:

$$G(r) = \frac{G_0}{\epsilon r + 1} \quad (3)$$

$$f(r) = 1 + 3G_0 M_0 \epsilon - \frac{2G_0 M_0}{r} - (1 + 6\epsilon G_0 M_0)\epsilon r - \frac{\Lambda_0 r^2}{3} + r^2 \epsilon^2 (6\epsilon G_0 M_0 + 1) \ln \left(\frac{c_4(\epsilon r + 1)}{r} \right) \quad (4)$$

$$\Lambda(r) = \frac{-72\epsilon^2 r(\epsilon r + 1) \left(\epsilon r + \frac{1}{2}\right) \left(G_0 M_0 \epsilon + \frac{1}{6}\right) \ln \left(\frac{c_4(\epsilon r + 1)}{r} \right) + 4r^3 \Lambda_0 \epsilon^2 + (12\epsilon^3 + 6\Lambda_0 \epsilon + 72\epsilon^4 G_0 M_0) r^2}{2r(\epsilon r + 1)^2} \quad (5)$$
$$+ \frac{(72\epsilon^3 G_0 M_0 + 11\epsilon^2 + 2\Lambda_0) r + 6\epsilon^2 G_0 M_0}{2r(\epsilon r + 1)^2}.$$

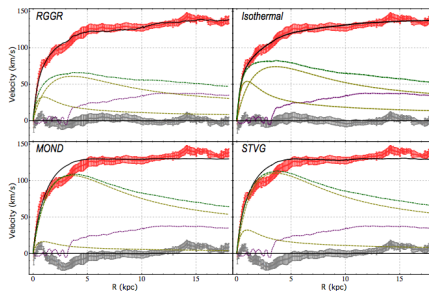
Constants of integration: G_0 , M_0 , Λ_0 , ϵ , c_4



Within gravity

Scale dependence

Effect on rotation curves



Comparison with MOND, DM, ...
promising



Within gravity

Scale dependence

Effect on lensing

... **not so promising**

Could be there but unlikely to explain it all

Try particles



Within gravity

Scale dependence

Effect on lensing

... **not so promising**

Could be there but unlikely to explain it all

Try particles



Particles



Particles

New particles, new game:



Anything goes as long as...



New particles, new game:

Anything goes as long as...

- Right masses (not too light, rotation curves; not too heavy abundance and observability)
- right couplings (abundance, stability)
- right “non-couplings” collider and indirect detection constraints
- ...
- Our contribution in this direction, Higgs sector*, Susy-gravitino sector**

*M.A. Diaz, B.K., [S. Urrutia-Quiroga](#) arXiv:1511.04429

**M.A. Diaz, [S. Garcia](#), B.K., Phys.Rev. D84 (2011) 055007



Particles

Higgs

Simple SM extension: Inert Higgs Doublet Model (IDM) by N. G. Deshpande y E. Ma en 1978

- In addition to the usual SM Higgs doublet (Φ_S) one introduces an additional doublet (Φ_D)
- Discrete symmetry \mathbb{Z}_2 such that

$$\begin{aligned}\Phi_S &\mapsto \Phi_S \\ \Phi_D &\mapsto -\Phi_D \\ SM &\mapsto SM\end{aligned}$$

- Thus Φ_D has no tree level couplings to SM fermions
- Phenomenology compatible with SM



Particles

Higgs

IDM potential:

$$V = \mu_1^2 \mathbf{A} + \mu_2^2 \mathbf{B} + \lambda_1 \mathbf{A}^2 + \lambda_2 \mathbf{B}^2 + \lambda_3 \mathbf{A} \mathbf{B} + \lambda_4 \mathbf{C} \mathbf{D} + \frac{\lambda_5}{2} (\mathbf{C}^2 + \mathbf{D}^2)$$

where $\mathbf{A}, \mathbf{B}, \mathbf{C}, \mathbf{D}$ are given by

$$\mathbf{A} = \Phi_S^\dagger \Phi_S, \quad \mathbf{B} = \Phi_D^\dagger \Phi_D, \quad \mathbf{C} = \mathbf{D}^\dagger = \Phi_S^\dagger \Phi_D$$

- Only Φ_S acquires VEV since want to preserve \mathbb{Z}_2 symmetry
- **Degrees of freedom:**
8 (two doublets) – 3 (Goldstone) = 5 (Physical scalars)
- **Parameters:**
7 (potential) – 2 (M_Z, M_H) = 5 (free parameters)



Particles

Higgs

Physical content:

- $\Phi_S : h$, *SM-like*, with Yukawa couplings
- $\Phi_D : H, A, D^\pm$, inert scalars with interior couplings and couplings to EW gauge bosons through kinetic term
- \mathbb{Z}_2 symmetry
 - Lightest inert scalar stable
 - **DM** candidate
- Parameters:

$$M_H, M_A, M_D \quad \text{y} \quad \lambda_2, \lambda_{345} \equiv \lambda_3 + \lambda_4 + \lambda_5$$



Particles

Higgs

Restrictions:

- 1 Positive potential, minimum, perturbativity and unitarity, inert vacuum
- 2 Several electroweak precision tests
- 3 Width of electroweak gauge bosons Z and W
- 4 DM candidate has to be H (neutral)
- 5 LHC restrictions: $\text{Br}(h \rightarrow \text{invisible}) < 0,43$

Now scan within those restrictions



Particles

Higgs

Parameter scan and **DM** density:

- Random scan over previously mentioned parameters within restrictions
- Mass range $1 \text{ GeV} \leq M_{H,A,D} \leq 1 \text{ TeV}$
- Check *relic density* (WMAP, Planck) $\Omega_{DM} h^2 = 0,1181 \pm 0,0012$
- Cosmological parameters obtained with micrOMEGAs
- Tolerance in relic density $\pm 3\sigma$

Color coding	
●	Relic density too low $< \rho_{DM} \pm 3\sigma$
●	Relic density within $\pm 3\sigma$
●	Relic density too high $> \rho_{DM} \pm 3\sigma$

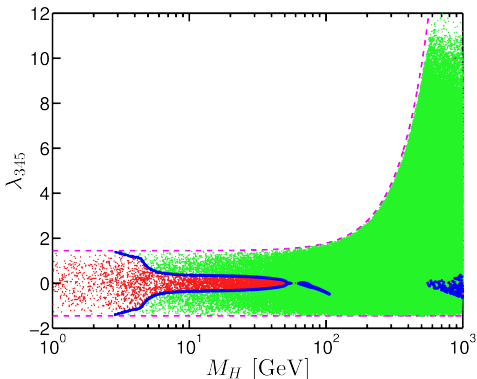


Particles

Higgs

Results scan:

Projection to (M_H, λ_{345}) plane



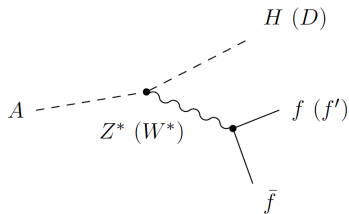
Note: Upper line from inert vacuum, lower line from vacuum stability condition



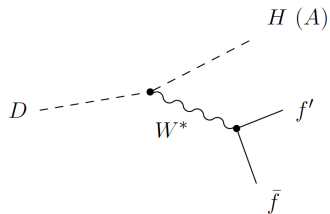
Particles

Higgs

For collider signals also study BR inert decays:



(a)



(b)

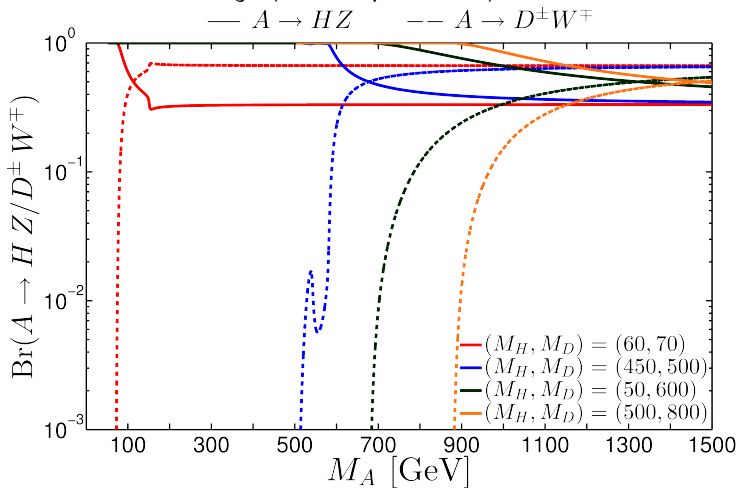
- Only depend on inert scalar masses + SM (No $\lambda...$)
- Take scalar masses on-shell



Particles

Higgs

Results inert decays (no λ dependence):



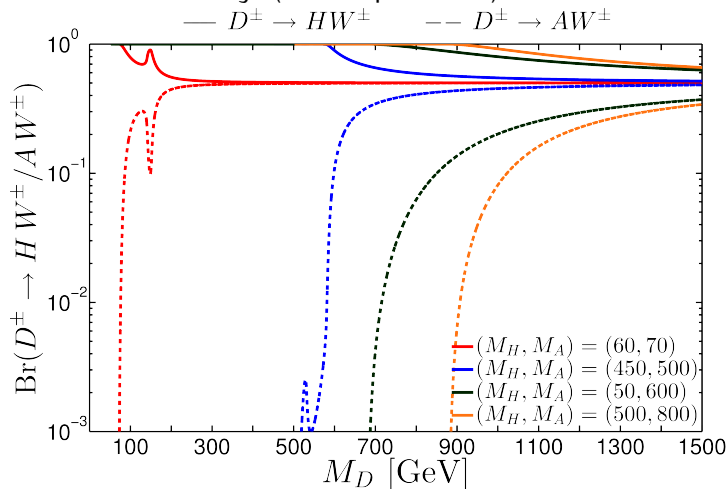
Crossover



Particles

Higgs

Results inert decays (no λ dependence):



No crossover, decay to DM H always larger

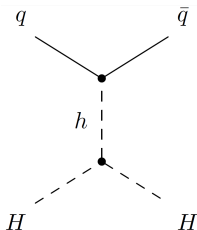


Particles

Higgs

Results direct detection:

- Scattering DM–nucleon en el IDM (“Higgs portal”)



$$\sigma^{SI}_{DM-N} = \frac{\lambda_{345}^2}{(4\pi M_h^4)} \frac{M_N^4 f_N^2}{(M_H + M_N)^2}$$

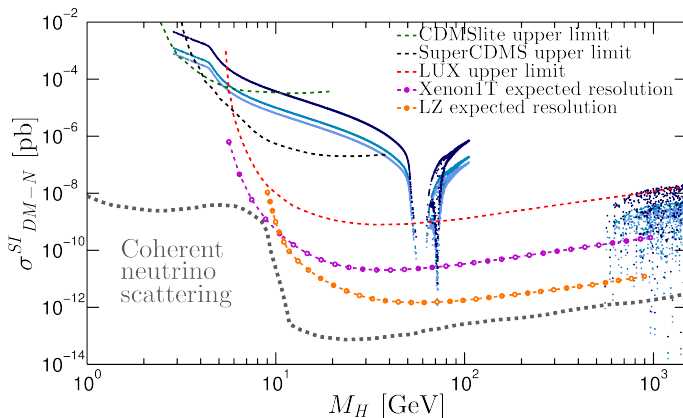
- f_N comes from QCD take conservative small values
- Tree level calculation
- Compare to upper limits from recent experiments



Particles

Higgs

Results direct detection (only take right **DM** density):



Drop when $m_H = 2m_h$ (efficient h production),
model will be largely testable



Light particles high energy signals

Light particles high energy signals



Light particles high energy signals

“High energy” signals

DM particles and “high energy” signals

- Collider production (just saw example)
- Annihilation work in progress
- Decay studied example*
- Acceleration of light DM particles**

*M. Diaz, S. Garcia, B.K. Phys.Rev. D84 (2011) 055007 **C. Armaza, M. Banados, B.K. arXiv:1510.01223



Light particles high energy signals

“High energy” signals

- Black holes can in principle produce $E_{CM} \rightarrow \infty$, but one needs
- **Extremely rotating** black hole
- Collision at the **horizon**
- Angular momentum **/: critical**

⇒ **Unlikely**, hard to observe



Light particles high energy signals

Idea:

Let the particle rotate and the black hole be spherical

- Can one produce $E_{CM} \rightarrow \infty$?

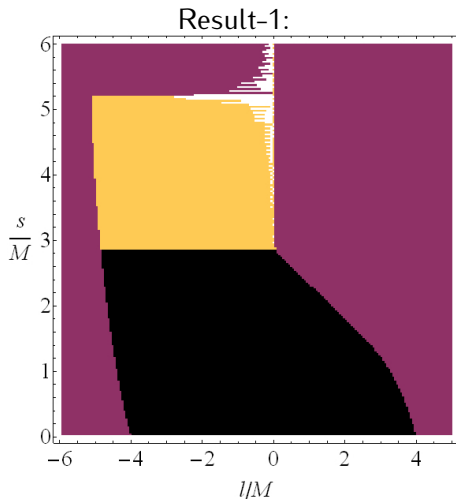
If yes:

- Has the collision to be at the **horizon**?
- Has the angular momentum **J : to be critical**?
- Is there a notion of **extremely rotating** particle?

⇒ Solve Papapetru equations (geodesics modified by spin) and see ...



Light particles high energy signals



E_{CM} divergent for yellow region



Particles

“High energy” signals

That was the **good news**,

what are the **bad news**?

- For solar mass object and spin 1/2 need $m_{DM} \approx 10^{-15}$ eV
- Papapetru equations allow for solutions with **superluminal** regime???



More Ideas?

Working on more ideas for the the dark matter sector?

Yes: Gamma rays from the galactic center ...

Juan Antonio Paredes and German Gomez



Take home message

DM is still mysterious,
but the good thing is that
a lot of observational evidence coming up!



Thank you

Thank you !

