

Vacuum energy, Casimir effect, and Newton's non-constant

B. Koch

with R. Sedmik, M. Pitschmann, and C. Käding



Nottingham
26.10.2022

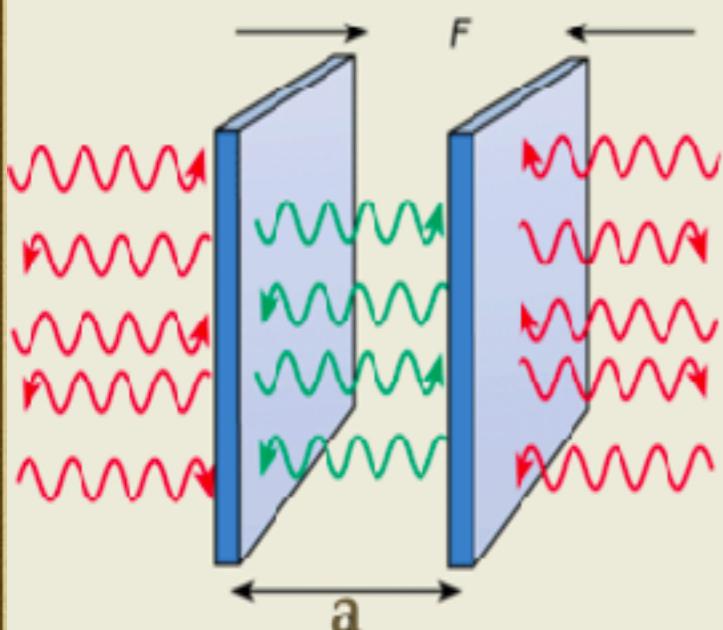
Based on
arXiv:SOON

Content

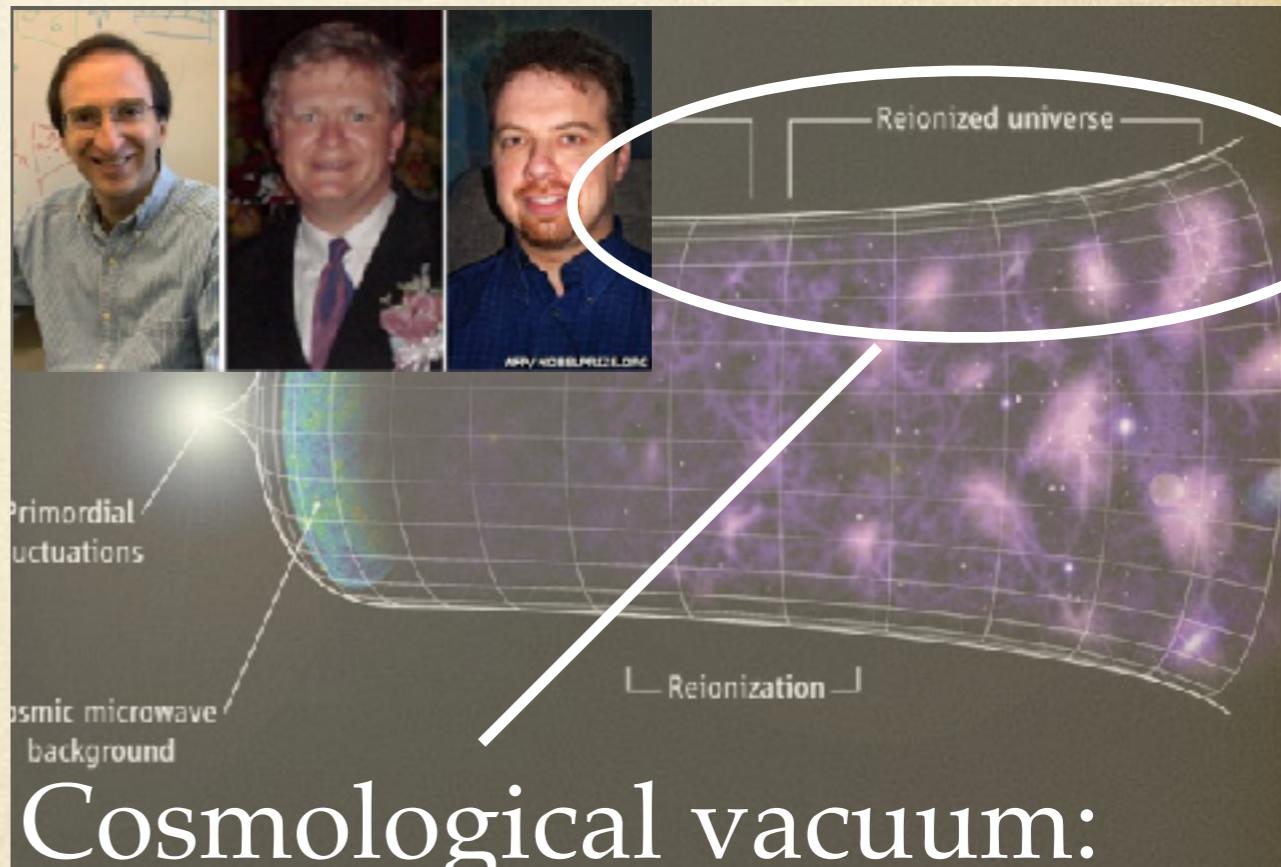
- What we would like to know
- Vacuum energy density - in the lab
- Vacuum energy density - in the Universe
- Vacuum energy density - combined
- Scale-dependent framework
- Towards experiment
- Discussion and Conclusion

Vacuum energy

What we know so far?



*



Quantum vacuum:

$$\rho_Q$$

accelerates plates

$$\rho_\Lambda$$

accelerates Universe

Vacuum energy

What we would like to know:

How are they related?

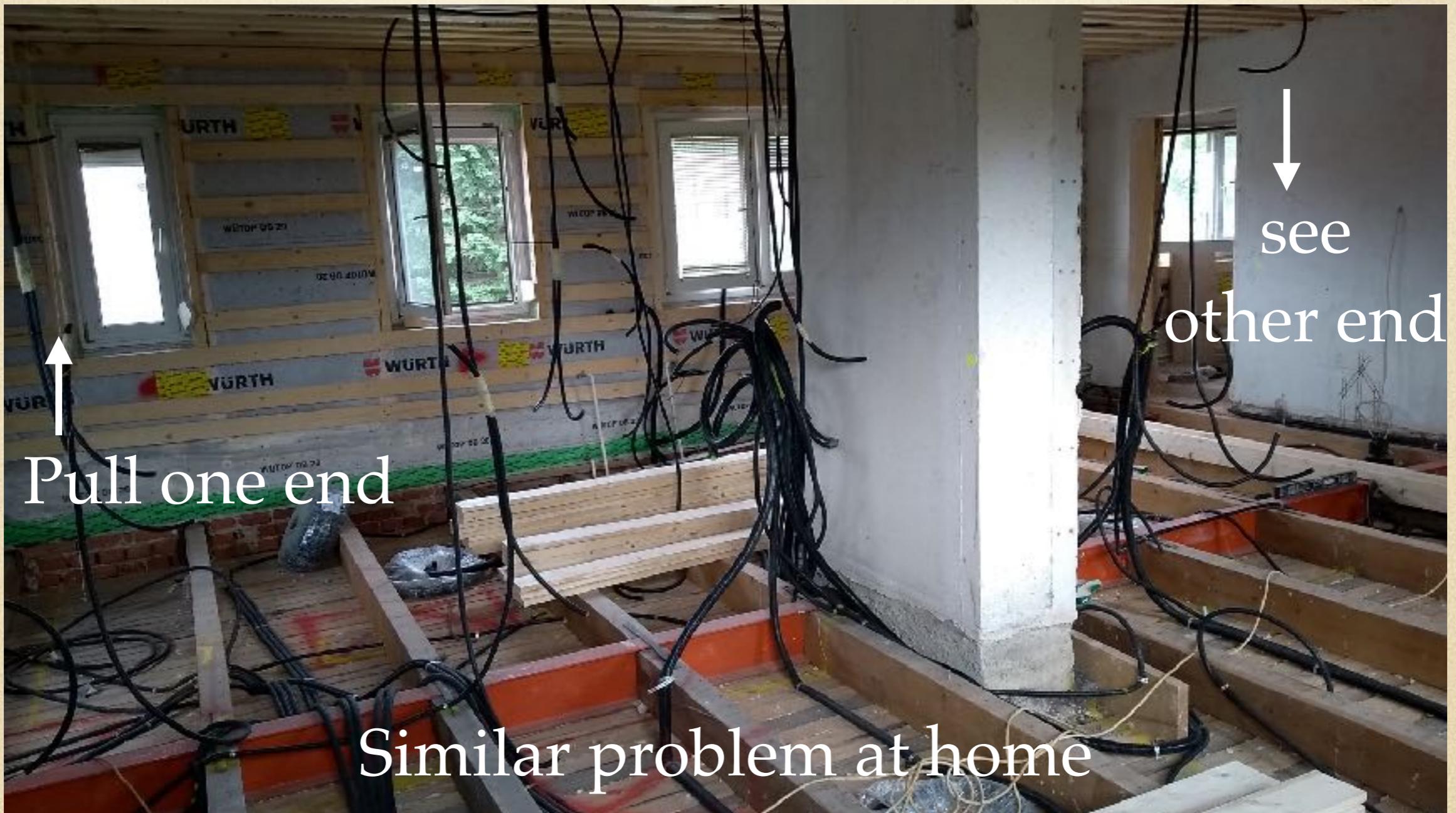
$$\rho_Q$$

$$\rho_\Lambda$$



Vacuum energy

How are they related?

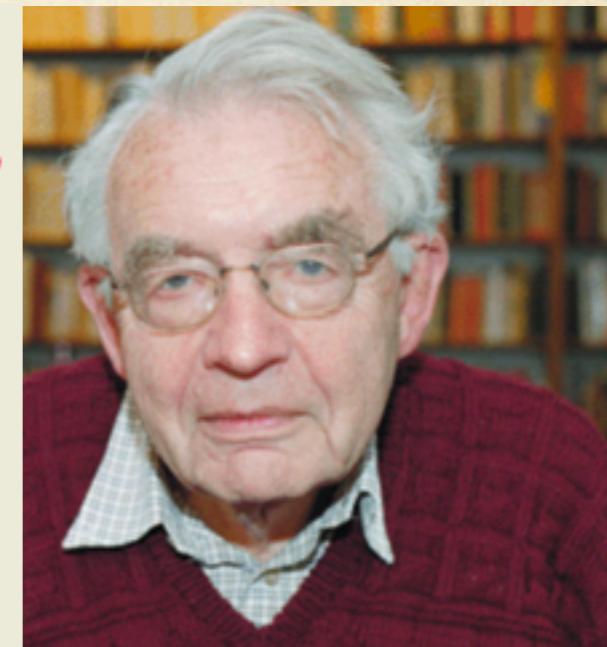
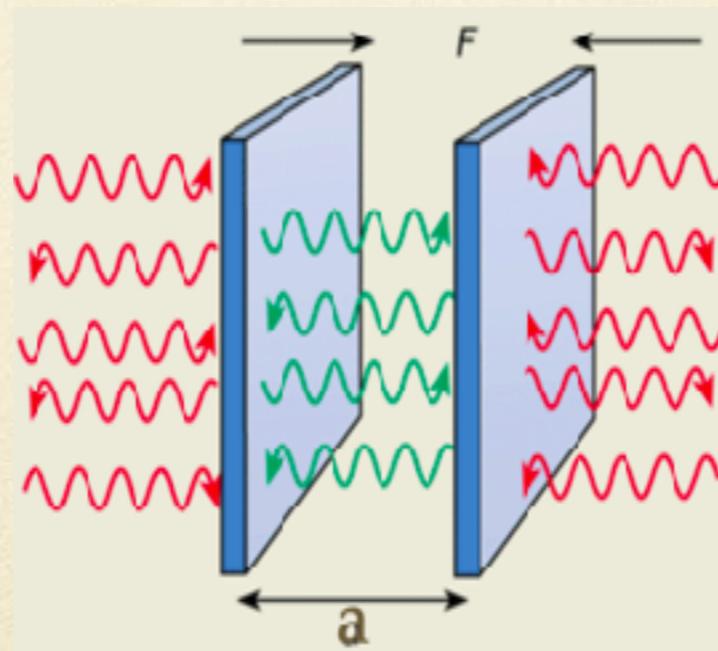


In the lab

Casimir effect

Predicted 1948

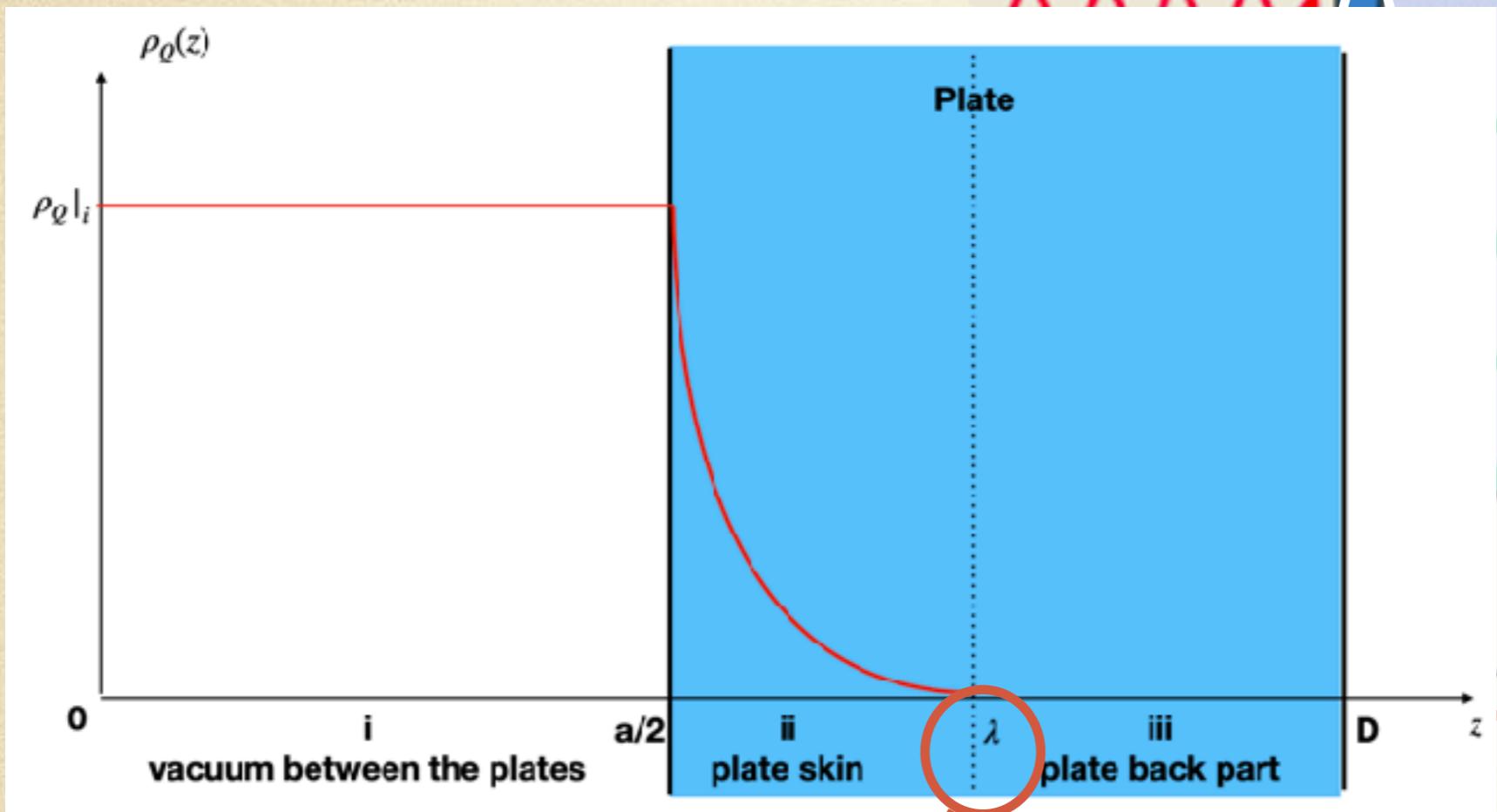
Observed 1997^{ref [10]}



$$\rho_C = -\frac{\hbar\pi^2}{720a^4}$$
$$\Rightarrow \frac{F_Q}{A} \approx \rho_Q \cdot a$$

In th.

In



$$\vec{F}_{G,12} = -\vec{F}_{G,21} = G_0 \int_{V_2} d^3x_2 \int_{V_1} d^3x_1 \frac{\rho_M(\vec{x}_1) \rho_M(\vec{x}_2) (\vec{x}_2 - \vec{x}_1)}{|\vec{x}_2 - \vec{x}_1|^3}$$

- Finite penetration depth

$$\lambda \approx 10^{-8} m,$$

thus $\rho_C = \rho_C(\vec{x})$

remember

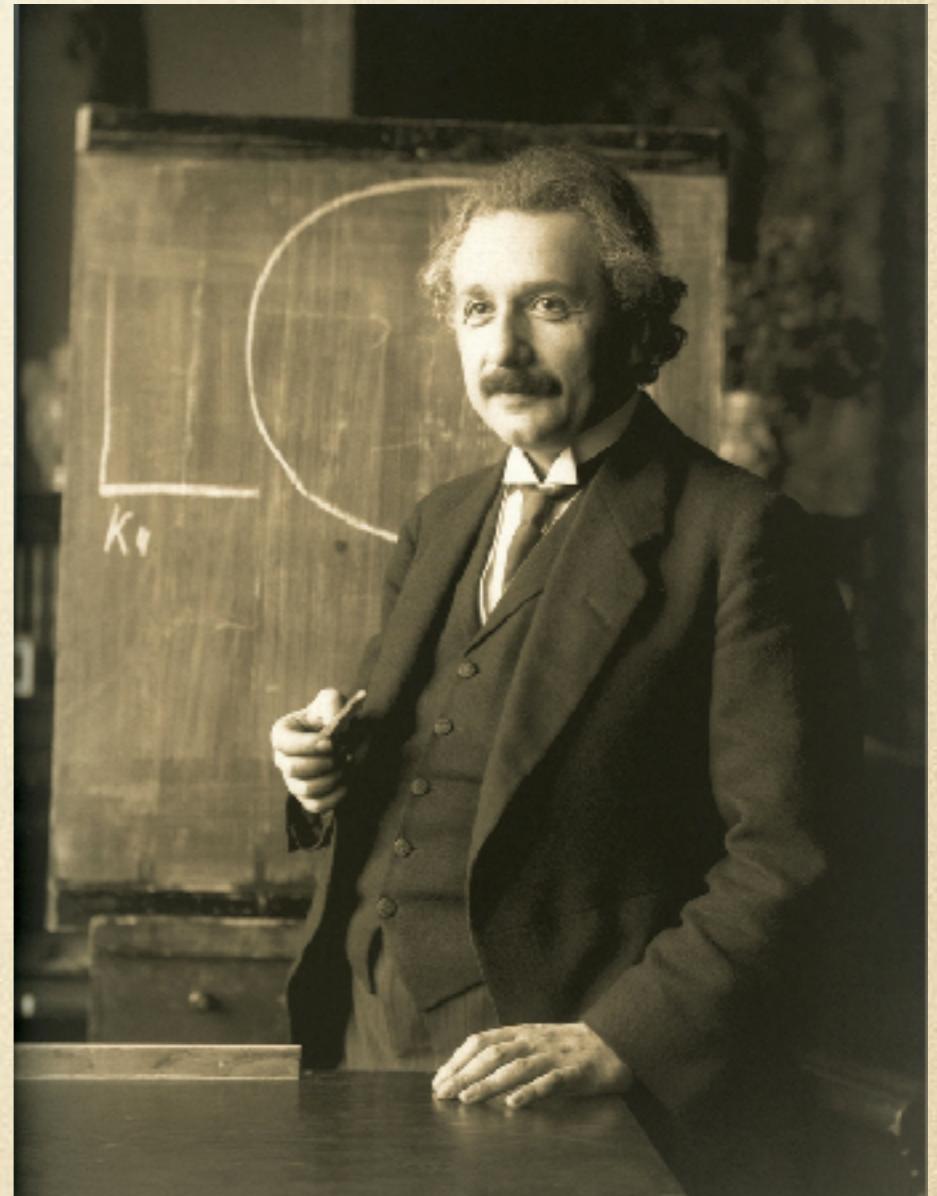
In the Universe

Albert Einstein

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

⇒ Friedman eq.

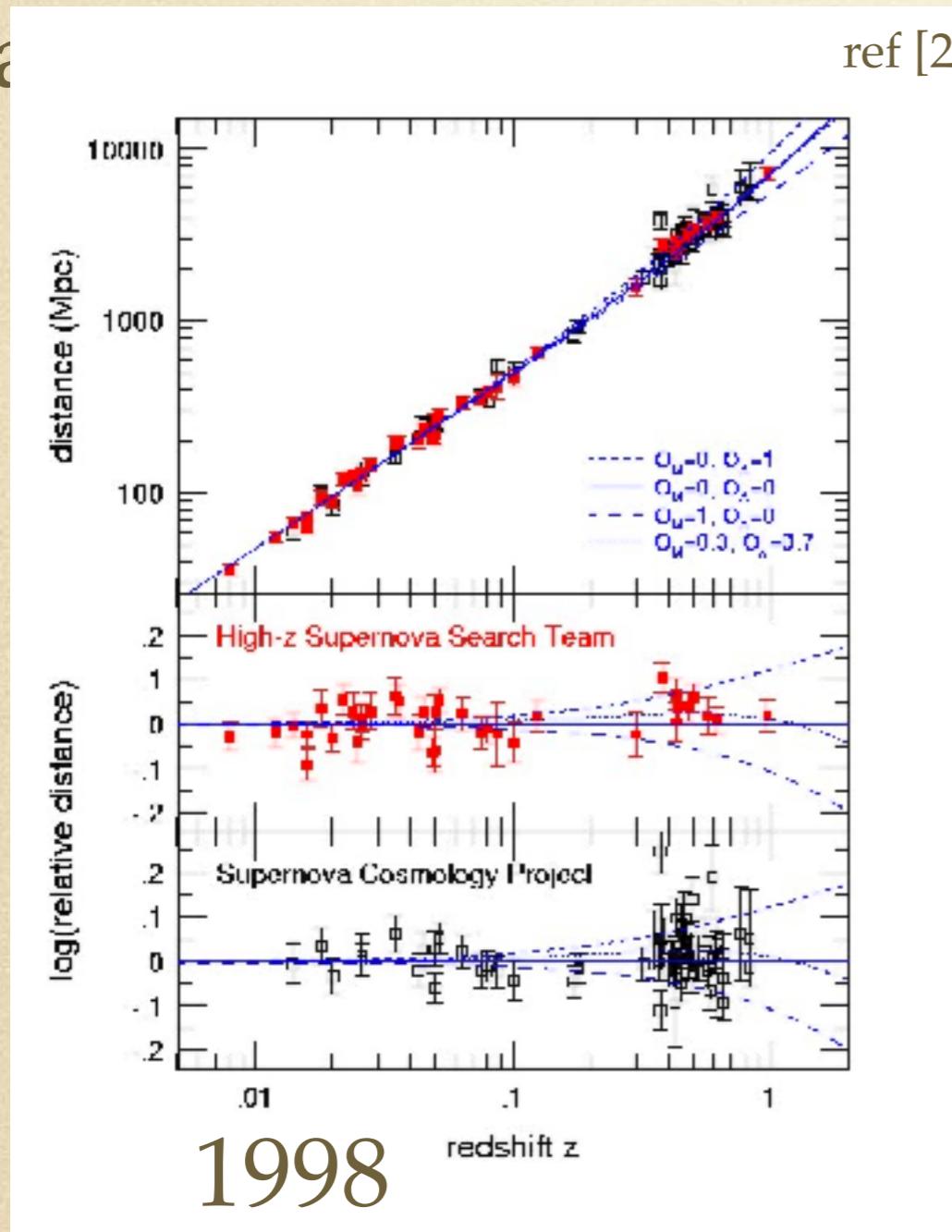
$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \frac{\Lambda}{3} \equiv 0$$



In the Universe

S. Perlmutter, A. Riess, B. Schmidt, & others

mea



NP 2011

$$\dot{a} \neq 0$$

$$\ddot{a} > 0 \Rightarrow \Lambda > 0 !$$

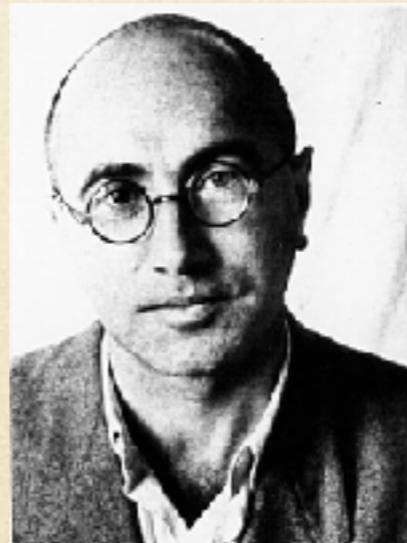
In the Universe

Λ as an energy density

$$\rho_{\Lambda,0} = \frac{\Lambda_0 c^4}{8\pi G_0} = 5.35 \times 10^{-10} J/m^3$$

Quantum origin?

Yakov Zeldovich, 1967



ref [4]
Big theoretical
puzzle

Steven Weinberg, 1998



In the Universe

Big theoretical puzzle

In short QFT with cutoff $\rho_Q \sim c\kappa_0^4/\hbar^3$

As ratio

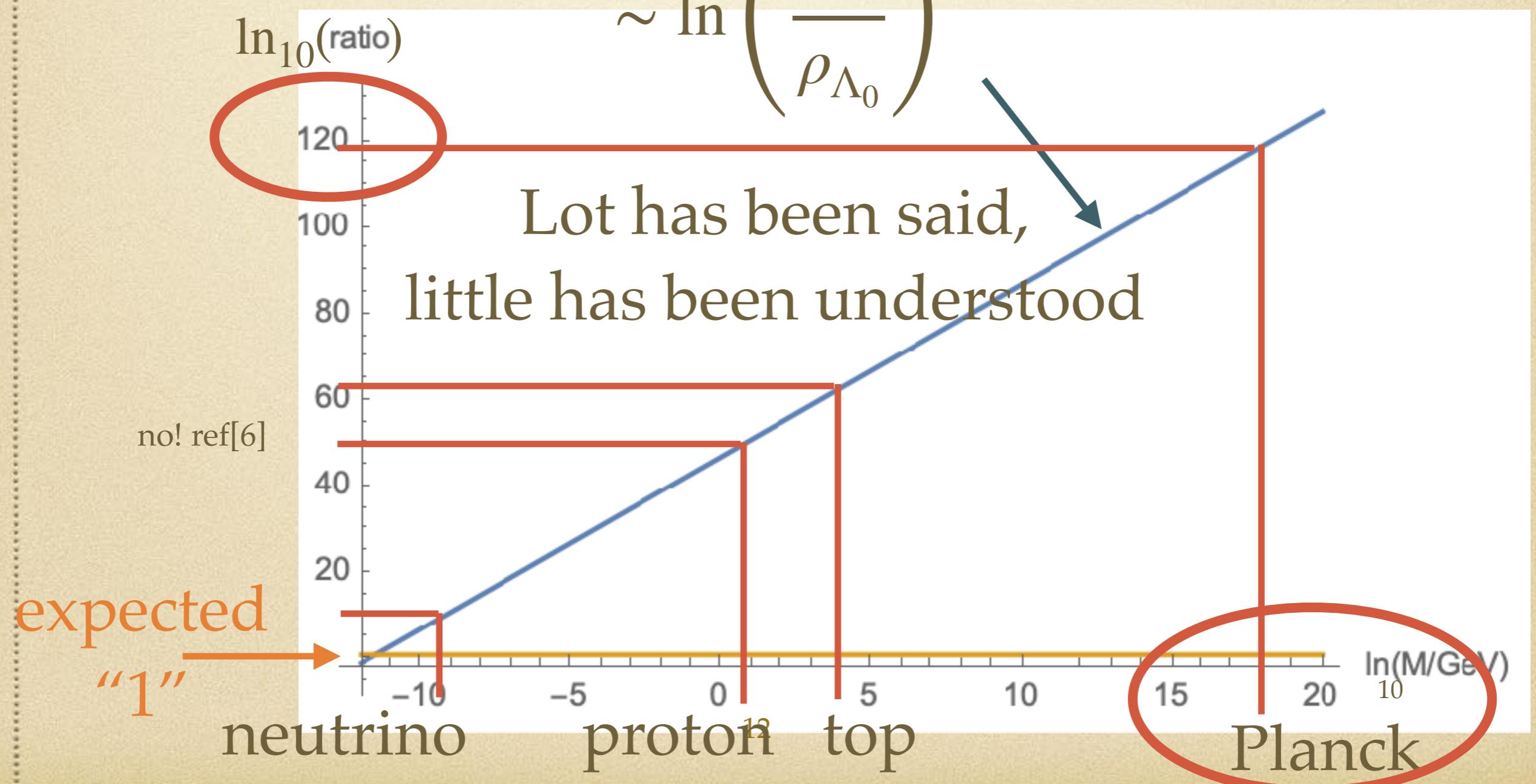
$$\Upsilon_0 \equiv \frac{\rho_{\Lambda_0}}{\rho_{Q,0}(\kappa)} = \frac{\Lambda_0 c^3 \hbar^3}{8\pi G_0 \kappa_0^4} = \begin{cases} 10^{-121} & \text{for } \kappa_0 = c \sqrt{\frac{c\hbar}{G_0}} \\ 10^{-55} & \text{for } \kappa_0 = cm_Z. \end{cases}$$

In the Universe

Big theoretical puzzle

Problem as a ratio:

$$\sim \ln \left(\frac{\kappa^4}{\rho_{\Lambda_0}} \right)$$



In the Universe

Big theoretical puzzle

experimental input needed

use Casimir experiment to gain insight

Change ρ_Q

See whether ρ_Λ

changes

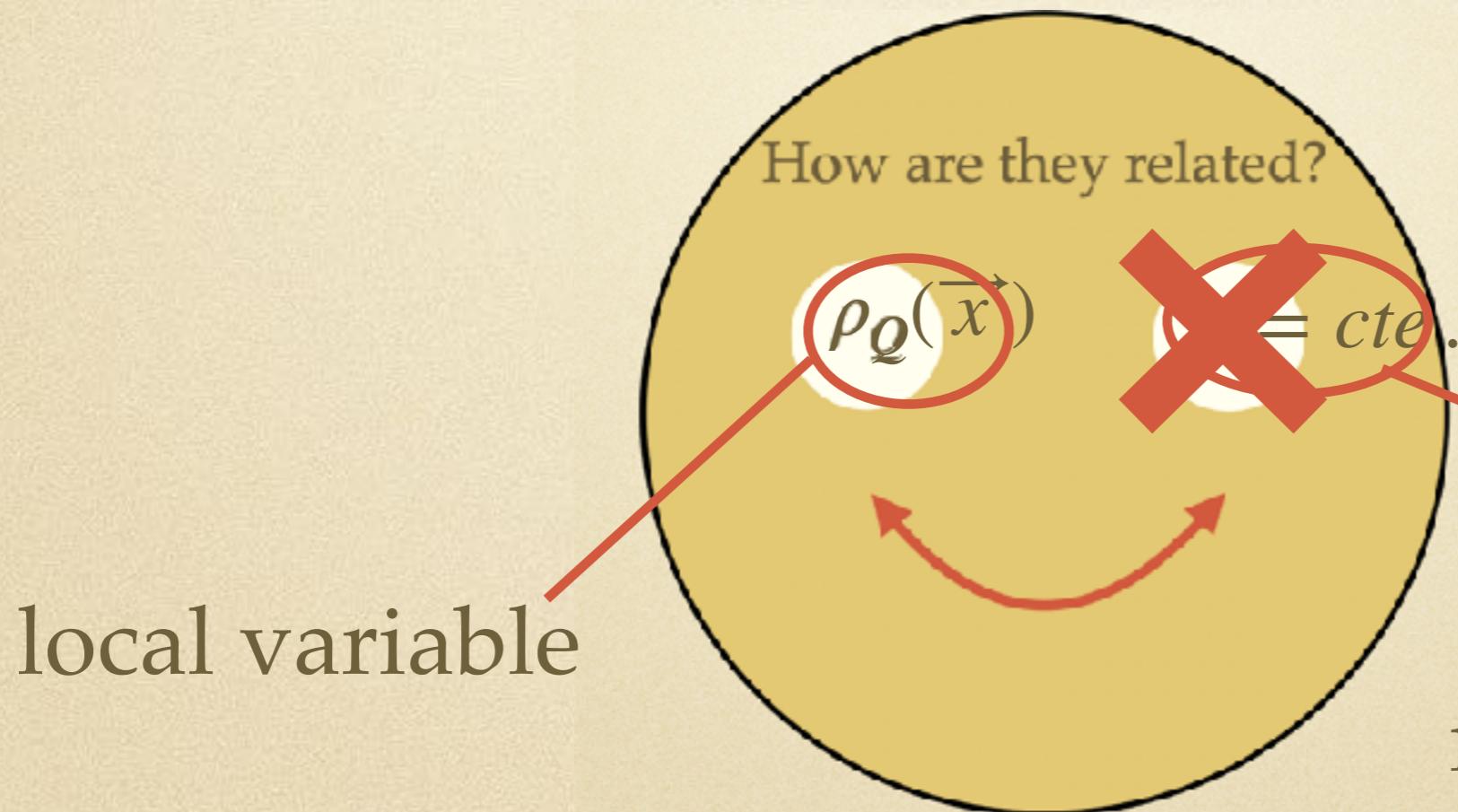


Hypothesis

Parametrize small change

original cosmo $\rightarrow \rho_{\Lambda_0} - \alpha \cdot \rho_C = \rho_\Lambda \leftarrow$ modified cosmo

quantum modification



wait ...

universal
constant

need variable,
 $\rho_\Lambda(\vec{x}), \Lambda(\vec{x}), \dots$

Hypothesis

...need variable, $\rho_{\Lambda}(\vec{x})$, $\Lambda(\vec{x})$, but first...

Opinion poll:

What will be found?

$$\rho_{\Lambda_0} - \alpha \cdot \rho_C = \rho_{\Lambda}$$

huge bounds

$\alpha < \text{huge}$

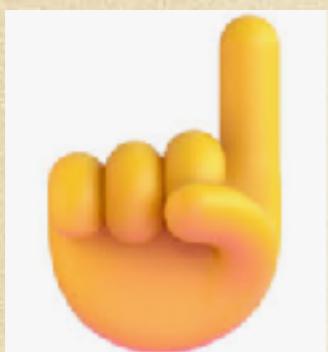


small bounds no bounds

discovery



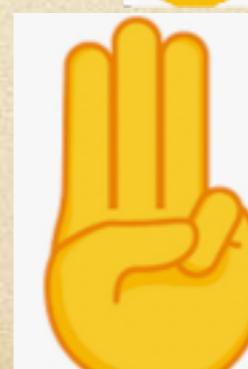
$\alpha \approx \text{some} \pm \Delta$



$\alpha < \text{small}$



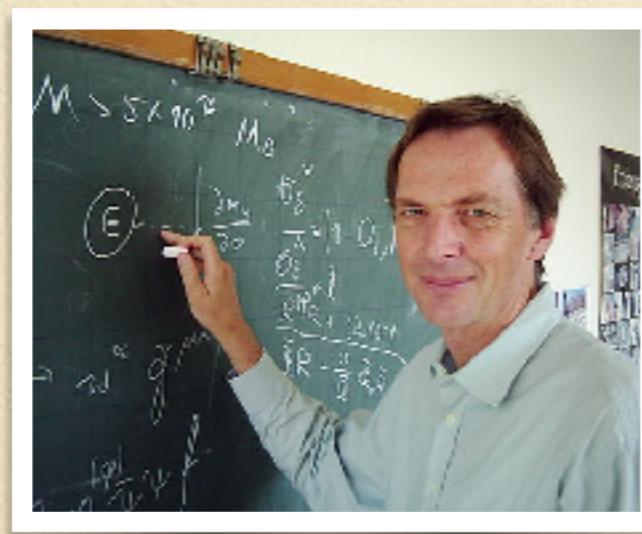
\emptyset



Scale Dependent Framework

Gravity as effective QFT

$$\Gamma_k = \int d^4x \sqrt{-g} \left(\frac{R}{G_k} - 2 \frac{\Lambda_k}{G_k} \right) + \dots$$



variable
couplings

Weinberg,

Wetterich, Donoghue,

SD-Framework

scale-dependent (SD) couplings $\Lambda = \Lambda(k)$, $G = G(k)$,
two main consequences:

- Modified Einstein equations

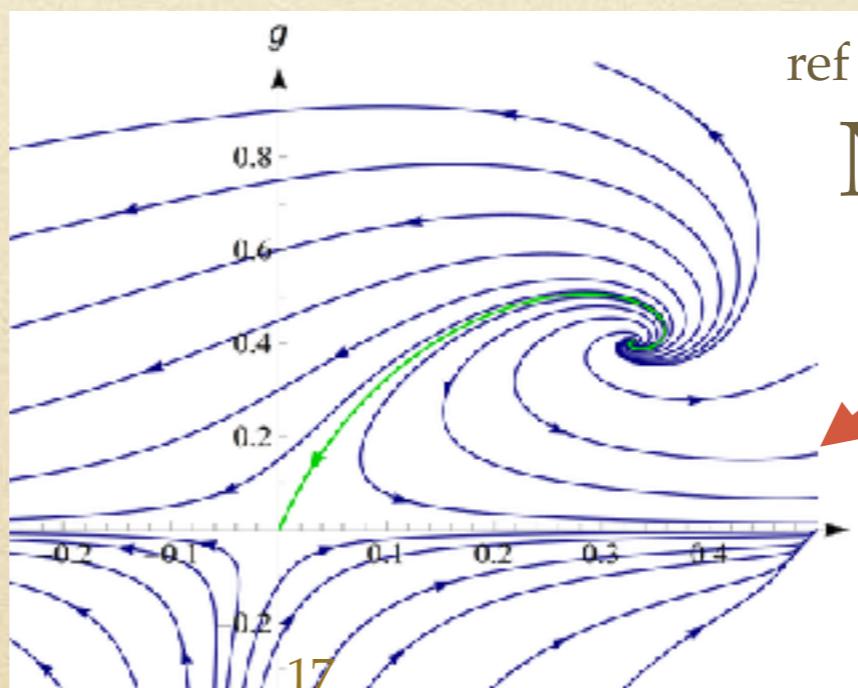
$$G_{\mu\nu} = 8\pi G(k)T_{\mu\nu} - \Lambda(k)g_{\mu\nu} - \Delta t_{\mu\nu}$$

- Couplings are connected in RG flow

example:

$$G_k = \frac{g_k}{k^2}$$

$$\Lambda_k = \lambda_k k^2$$



Not independent
both change
or
none changes

SD-Framework

- Couplings are connected in RG flow

We are only interested in SD small modifications
of the couplings in the IR:

Expand:

$$G(k) = G_0(1 + g(k)) = G_0 \left(1 + C_1 G_0 k^2 \right) + \mathcal{O}(k^4)$$

$$\Lambda(k) = \Lambda_0(1 + \lambda(k)) = \Lambda_0 \left(1 + C_3 G_0 k^2 \right) + \mathcal{O}(k^4)$$

Theorist:
predict

Phenomenologist:
use to predict

Experimentalist:
measure

SD-Framework

- Couplings are connected in RG flow

We are only interested in SD small modifications
of the couplings in the IR:

Expand:

$$G(k) = G_0(1 + g(k)) = G_0 \left(1 + C_1 G_0 k^2 \right) + \mathcal{O}(k^4)$$

$$\Lambda(k) = \Lambda_0(1 + \lambda(k)) = \Lambda_0 \left(1 + C_3 G_0 k^2 \right) + \mathcal{O}(k^4)$$

In density

$$\rho_\Lambda(k) = \frac{c^4 \Lambda(k)}{8\pi G(k)} = \rho_0 - \cancel{k^2} c^4 \frac{(C_1 - C_3)}{8\pi} \Lambda_0 + \mathcal{O}(k^4)$$

?

SD-Framework

To identify k^2 compare

$$\rho_\Lambda(k) = \rho_{\Lambda_0} - k^2 c^4 \frac{(C_1 - C_3)}{8\pi} \Lambda_0$$

$$\rho_\Lambda = \rho_{\Lambda_0} - \alpha \cdot \rho_C(\vec{x})$$
$$\Rightarrow k^2(\vec{x}) = \alpha \frac{8\pi \rho_C(\vec{x})}{c^4(C_1 - C_3)\Lambda_0}$$

Insert in couplings

$$G(k(\vec{x})) = G_0 \left(1 + C_1 G_0 \alpha \frac{8\pi \rho_C(\vec{x})}{c^4(C_1 - C_3)\Lambda_0} \right)$$

Thus, ρ_Q modifies Newton coupling!

SD-Framework

- Modified Einstein equations

$$G_{\mu\nu} = 8\pi G(k)T_{\mu\nu} - \Lambda(k)g_{\mu\nu} - \Delta t_{\mu\nu}$$

Derivatives of
 $G(k(\vec{x}))$

Does Casimir/SD affect Newton?

Trace trick

$$R^\mu_\nu = 8\pi \frac{G(k)}{c^4} \left(T^\mu_\nu - \frac{1}{2} g^\mu_\nu T \right) + \Lambda(k)g^\mu_\nu + G(k) \left(\frac{1}{2} g^\mu_\nu \nabla^2 + \nabla^\mu \nabla_\nu \right) \frac{1}{G(k)}.$$

With non-relativistic matter

$$\left(T^\mu_\nu - \frac{1}{2} g^\mu_\nu T \right) = \frac{\rho_M}{2} \text{diag}(-1,1,1,1)$$

SD-Framework

- Modified Einstein equations

$$G_{\mu\nu} = 8\pi G(k)T_{\mu\nu} - \Lambda(k)g_{\mu\nu} - \Delta t_{\mu\nu}$$

Derivatives of
 $G(k(\vec{x}))$

Weak field and weak SD expansion...

$$ds^2 = -(1 + 2\epsilon_\Phi \Phi(r, \theta, \phi))c^2 dt^2 + (1 - 2\epsilon_\Phi \Psi(r, \theta, \phi))dr^2 + (1 + 2\epsilon_\Phi \Xi(r, \theta, \phi))r^2 d\Omega^2 + \mathcal{O}(\epsilon_\Phi^2)$$

$$G(k) = \epsilon_\Phi (G_0 + \epsilon_G \Delta G(k) + \mathcal{O}(\epsilon_\Phi^2))$$

$$\Lambda(k) \rightarrow \epsilon_\Phi \Lambda(k)$$

R_0^0 component

$$\vec{\nabla}^2 \Phi(r, \theta, \phi) = \frac{4\pi}{c^4} G_0 \rho_M(r, \theta, \phi) + \frac{\epsilon_G}{\epsilon_\Phi} \frac{\vec{\nabla}^2 \Delta G(k)}{2G_0} - \Lambda(k) + \mathcal{O}(\epsilon_\Phi, \epsilon_G)$$

Solution...

SD-Framework

- Modified Einstein equations

$$G_{\mu\nu} = 8\pi G(k)T_{\mu\nu} - \Lambda(k)g_{\mu\nu} - \Delta t_{\mu\nu}$$

Derivatives of
 $G(k(\vec{x}))$

Solution

$$\Phi(\vec{x}) = \frac{G_0}{c^4} \int_{V_1} d^3x_1 \frac{\tilde{\rho}_M(\vec{x}_1)}{|\vec{x} - \vec{x}_1|} + \mathcal{O}(\epsilon_\Phi)$$

with

$$\tilde{\rho}_M = \rho_m + c^2 \frac{\vec{\nabla}^2 G(k)}{8\pi G_0^2}$$

$$\alpha c^2 C_1 \frac{\vec{\nabla}^2 \rho_C(\vec{x})}{c^4(C_1 - C_3)\Lambda_0}$$

For extended objects...

SD-Framework

- Modified Einstein equations

$$G_{\mu\nu} = 8\pi G(k)T_{\mu\nu} - \Lambda(k)g_{\mu\nu} - \Delta t_{\mu\nu}$$

Derivatives of
 $G(k(\vec{x}))$

For extended objects: Newton

$$\vec{\mathcal{F}}_{G,12} = -\vec{\mathcal{F}}_{G,21} = G_0 \int_{V_2} d^3x_2 \int_{V_1} d^3x_1 \frac{\tilde{\rho}_M(\vec{x}_1)\tilde{\rho}_M(\vec{x}_2)(\vec{x}_2 - \vec{x}_1)}{|\vec{x}_2 - \vec{x}_1|^3}$$

with

$$\tilde{\rho}_M = \rho_m + c^2 \frac{\vec{\nabla}^2 G(k)}{8\pi G_0^2}$$

$$\alpha c^2 C_1 \frac{\vec{\nabla}^2 \rho_C(\vec{x})}{c^4(C_1 - C_3)\Lambda_0}$$

SD-Framework

- Modified Einstein equations &
- Couplings are connected in RG flow

⇒ Gravitational attraction between plates changes

$$\alpha c^2 C_1 \frac{\vec{\nabla}^2 \rho_C(\vec{x})}{c^4 (C_1 - C_3) \Lambda_0} \rightarrow \vec{\mathcal{F}}_{G,12} \neq \vec{F}_{G,12}$$

Hypothesis
can be tested by experiment:
Verify, or set bounds on α ...

SD-Framework

Gravitational attraction between plates changes

$$\alpha c^2 C_1 \frac{\vec{\nabla}^2 \rho_C(\vec{x})}{c^4(C_1 - C_3)\Lambda_g} \rightarrow \vec{\mathcal{F}}_{G,12} \neq \vec{F}_{G,12}$$

What that means

Change ρ_Q

Hypothesis

ρ_Λ not measurable

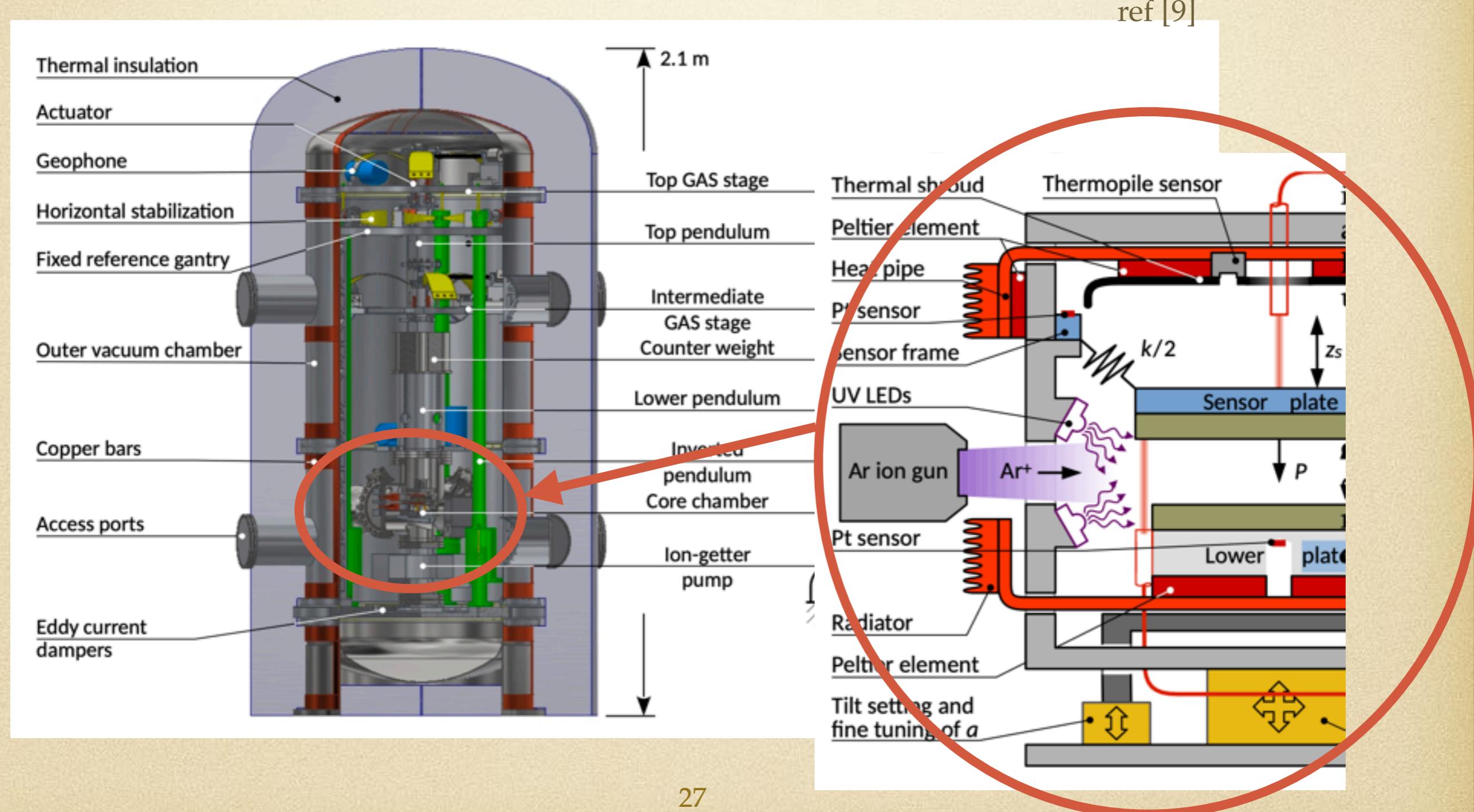


SD & flow

$\rho_M \rightarrow \tilde{\rho}_M$
measurable

Towards experiment

Cannex planned experiment



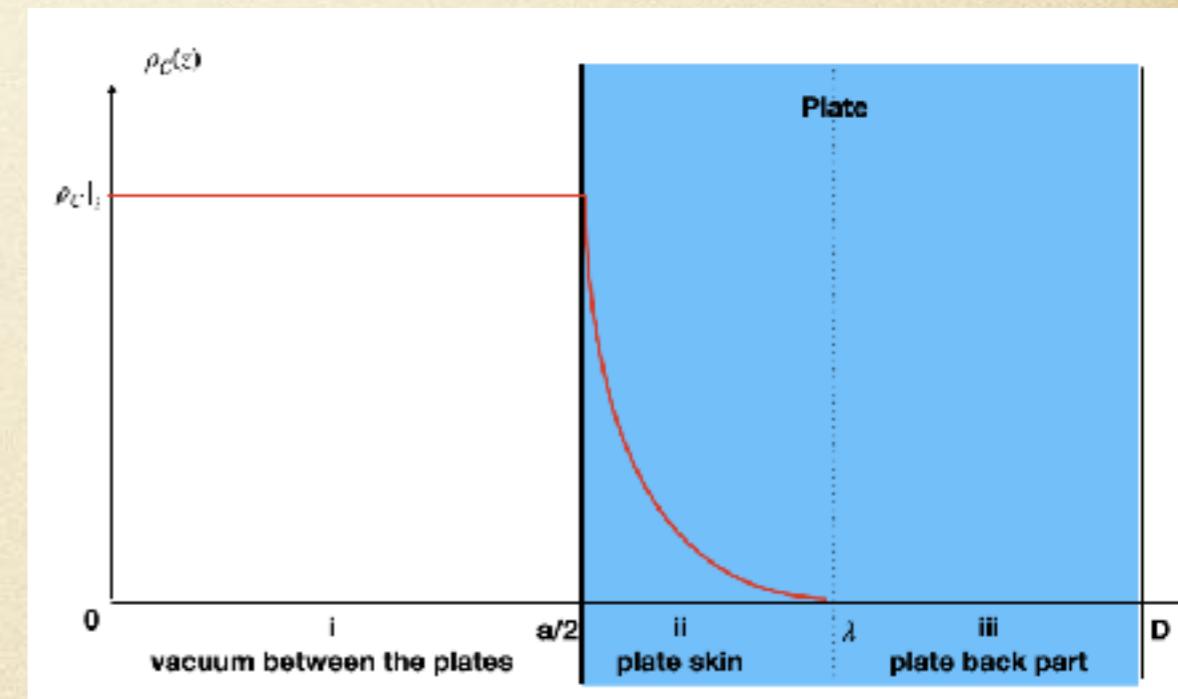
Towards experiment

Results (preliminary estimate):

$$1 \ll (8\pi G_0^2) \frac{\int_{a/2}^D dz \rho_M(z)}{\int_{a/2}^D dz c^2 \vec{\nabla}^2 G}$$

$$\alpha \frac{C_1}{C_1 - C_3} \ll 10^{-32}$$

$\alpha < small$



heading towards
ultra small

²⁸ potentially excludes models

Interpretation

$$\alpha \frac{C_1}{C_1 - C_3} \ll 10^{-32}$$

- A. ρ_Q contribution to ρ_Λ strongly suppressed ($\alpha \ll 1$)
- B. $\Lambda(k)$ has very weak RG coupling to $G(k)$
- C. Effective Einstein equations have additional fields, contributions, stuff, leading to cancellations...

For each interpretation many possible subcategories, e.g.

- B.
 - 1. Λ is not a coupling but a field
 - 2. G is not a coupling but a field
 - 3. RG group is not universal
 - 4. Hierarchy in QG parameters: $C_3 \gg C_1$
 - 5. ...

Interpretation

A: ($\alpha \ll 1$)

Implications for the CCP

$$\Upsilon_0 \equiv \frac{\rho_{\Lambda_0}}{\rho_{Q,0}(\kappa)} = \frac{\Lambda_0 c^3 \hbar^3}{8\pi G_0 \kappa_0^4} = \begin{cases} 10^{-121} & \text{for } \kappa_0 = c \sqrt{\frac{c\hbar}{G_0}} \\ 10^{-55} & \text{for } \kappa_0 = cm_Z. \end{cases}$$

Problem comes from the ambition

$$\rho_{\Lambda} = \Upsilon(\rho_Q) \cdot \rho_Q,$$

Casimir can contribute to both

$$\rho_Q = \rho_{Q,0} + \beta \cdot \rho_C$$

$$\rho_{\Lambda} = \rho_{\Lambda_0} - \alpha \cdot \rho_C$$

Should be $\beta = 1$,
but who knows ...

hypothesis,
 α

Interpretation

A: ($\alpha \ll 1$)

Implications for the CCP

Look at changes of the CCP

Find

$$\Upsilon'_0 \equiv \left. \frac{d\Upsilon(\rho_Q)}{d\rho_C} \right|_{\rho_C=0}$$

Measuring α

$$\alpha = \Upsilon'_0 + \beta \Upsilon_0 \quad \Rightarrow$$

Measure changes in CCP

Under construction



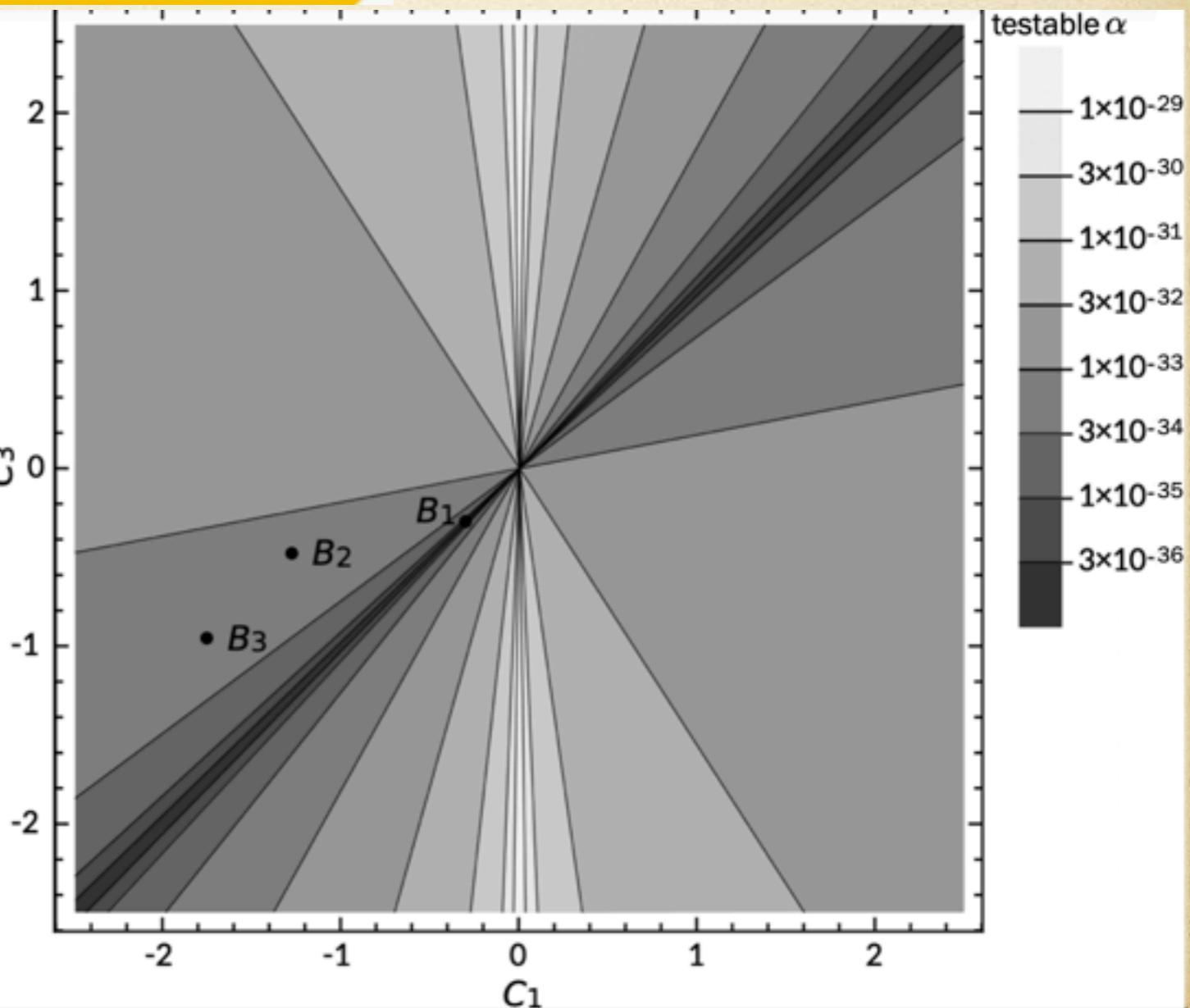
- Comparison with qu

	B_1	B_2	B_3
N_S	0	0	4
N_D	8	9	12
N_V	0	1	12
C_1	$-15/(16\pi)$	$-4/\pi$	$-11/(2\pi)$
C_3	$-15/(16\pi)$	$-3/(2\pi)$	$-3/\pi$
$\frac{C_1 K(C_1 - C_3)}{C_1 + C_3}$	∞	1.6	2.2

- More realistic simul
- Simulation for existi

- Implications for the

- ...



Some References

- 0) F. Canales, B. Koch, C. Laporte, A. Rincon JCAP no1 26, 2020
- 1) E. Hubble, Proceedings of the National Academy of Sciences of USA, Volume 15, Issue 3, pp. 168-173
- 2) Supernova Search Team (A. G. Riess (UC, Berkeley, Astron. Dept.) et al.). May 1998. 36 pp Published in Astron.J. 116 (1998) 1009-1038
- 3) R. J. Adler, B. Casey and O. C. Jacob, Am. J. Phys. 63, 620 (1995);
H. Martel, P. R. Shapiro and S. Weinberg, Astrophys. J. 492, 29 (1998);
S. Weinberg, Rev. Mod. Phys. 61, 1 (1989);
J. Martin Comptes Rendus Physique 13 (2012) 566-665.
- 4) M. Reuter, F. Saueressig, Phys.Rev. D65 (2002) 065016; and many others
- 5) e.g. M. Niedermaier and M. Reuter, Living Rev. Rel. 9, 5 (2006)
- 6) S.J. Brodsky and R. Shrock, Proc. Nat. Acad.Sci. 108, 45 (2011)
- 7) H. Fritzsch, Nucl. Phys. Proc. Suppl. 203-204, 3 (2010)
- 8) B.K., P. Rioseco, Carlos Contreras, PRD 91 (2015) no2, 025009.
- 9) R. Sedmik, M.Pitschann, Universe, 7, 234, (2021)
- 10) G. Bimonte, B. Spreng, P. A. Maia Neto, G.-L. Ingold, G. L. Klimchitskaya, V. M. Mostepanenko, and R. S. Decca, Universe 7, 93 (2021).
B. V. Derjaguin, I. I. Abrikosova, and E. M. Lifshitz, Q. Rev. Chem. Soc. **10**, 295 (1956).
S. Lamoreaux, Phys. Rev. Lett. **78**, 5 (1997).
U. Mohideen and A. Roy, Phys. Rev. Lett. **81**, 4549 (1998).
A. Roy, C.-Y. Lin, and U. Mohideen, Phys. Rev. D **60**, 111101 (1999).
G. Bressi, G. Carugno, R. Onofrio, and G. Ruoso, Phys. Rev. Lett. **88**, 041804 (2002).
R. S. Decca, D. López, B. Fischbach, and D. E. Krause, Phys. Rev. Lett. **91**, 050402 (2003).
J. M. Obrecht, R. J. Wild, M. Antezza, L. P. Pitaevskii, S. Stringari, and E. A. Cornell, Physical Review Letters **98**, 063201 (2007).
H. B. Chan, Y. Bao, J. Zou, R. A. Cirelli, F. Klemens, W. M. Mansfield, and C. S. Pai, Physical Review Letters **101**, 030401 (2008).
W. J. Kim, A. O. Sushkov, D. A. R. Dalvit, and S. K. Lamoreaux, Phys. Rev. Lett. **103**, 060401 (2009).
P. Zuurbier, S. de Man, G. Gruca, K. Heeck, and D. Iannuzzi, New Journal of Physics **13**, 023027 (2011), ISSN 1367-2630.
R. H. Schaefer, Ph.D. thesis, UC Riverside (2020).

Thank You!



Backup stuff

Scale-Dependent & FRG

Scale Dependent Framework

Gravity as effective QFT

$$\Gamma_{\textcolor{red}{k}} = \int d^4x \sqrt{-g} \left(\frac{R}{G_k} - 2 \frac{\Lambda_{\textcolor{red}{k}}}{G_k} \right) + \dots$$

Non renormalizable?

Yes, but ...

Could still be predictive QFT
(Asymptotic Safety)



Asymptotic Safety in a nutshell

$$\Gamma_{\textcolor{red}{k}} = \int d^4x \sqrt{-g} \left(\frac{R}{G_{\textcolor{red}{k}}} - 2 \frac{\Lambda_{\textcolor{red}{k}}}{G_{\textcolor{red}{k}}} \right) + \dots$$

- Idea: works if non trivial UV-fixed points for finite number of couplings (S.W)



Asymptotic Safety in a nutshell

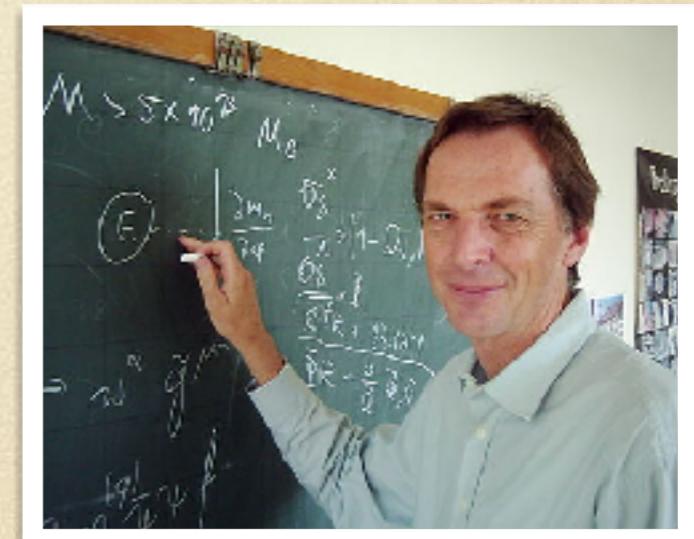
$$\Gamma_k = \int d^4x \sqrt{-g} \left(\frac{R}{G_k} - 2 \frac{\Lambda_k}{G_k} \right) + \dots$$

- Tool: Functional renormalization group equation

$$\partial_k \Gamma_k = \frac{1}{2} Tr \left(\frac{\partial_k R_k}{\Gamma_k^{(2)} + R_k} \right)$$

two point function

regulator



C. Wetterich

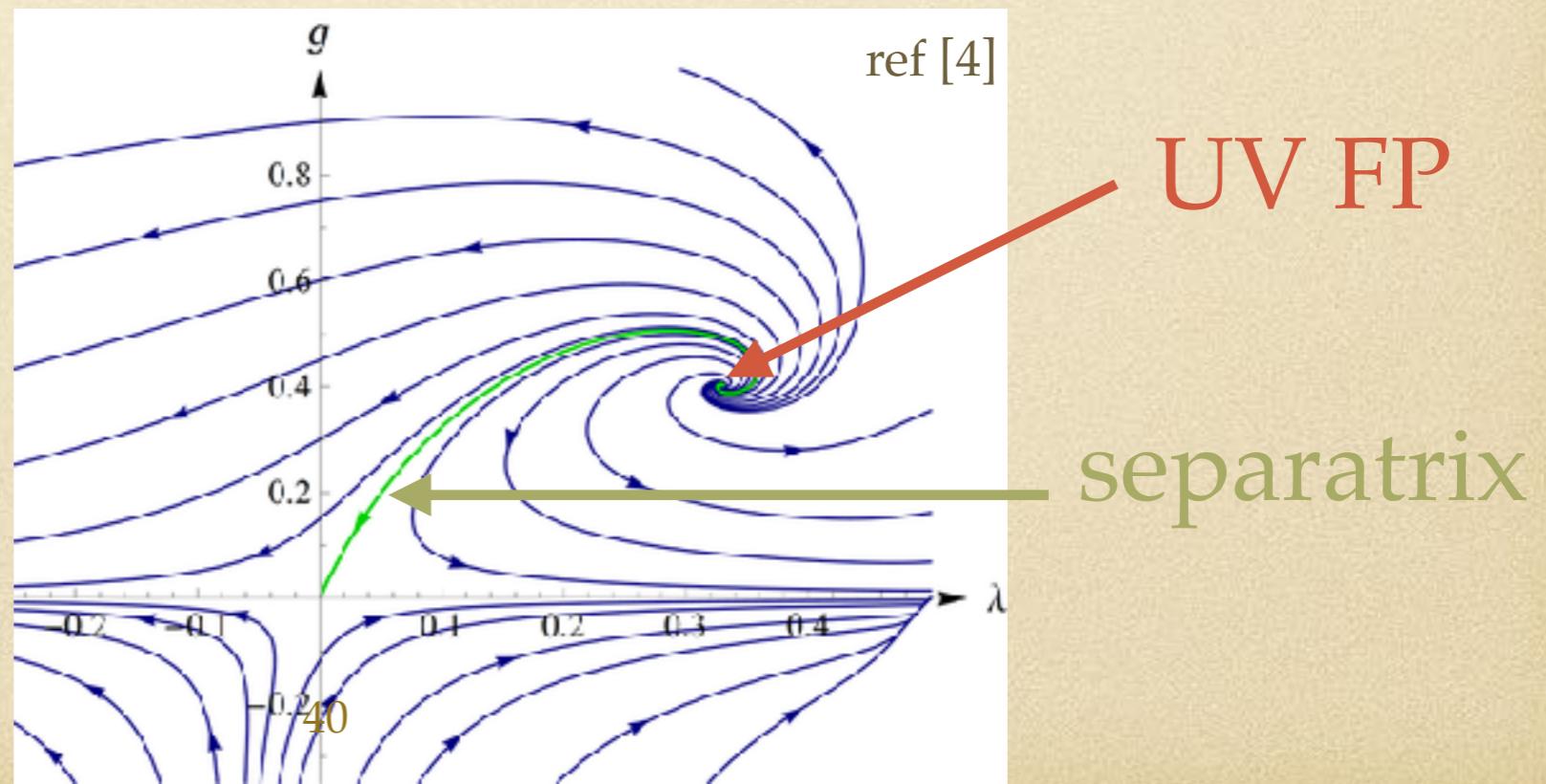
Asymptotic Safety in a nutshell

$$\Gamma_{\textcolor{red}{k}} = \int d^4x \sqrt{-g} \left(\frac{R}{G_{\textcolor{red}{k}}} - 2 \frac{\Lambda_k}{G_{\textcolor{red}{k}}} \right) + \dots$$

- Results: Plenty of evidence supporting idea

$$G_k = \frac{g_k}{k^2}$$

$$\Lambda_k = \lambda_k k^2$$

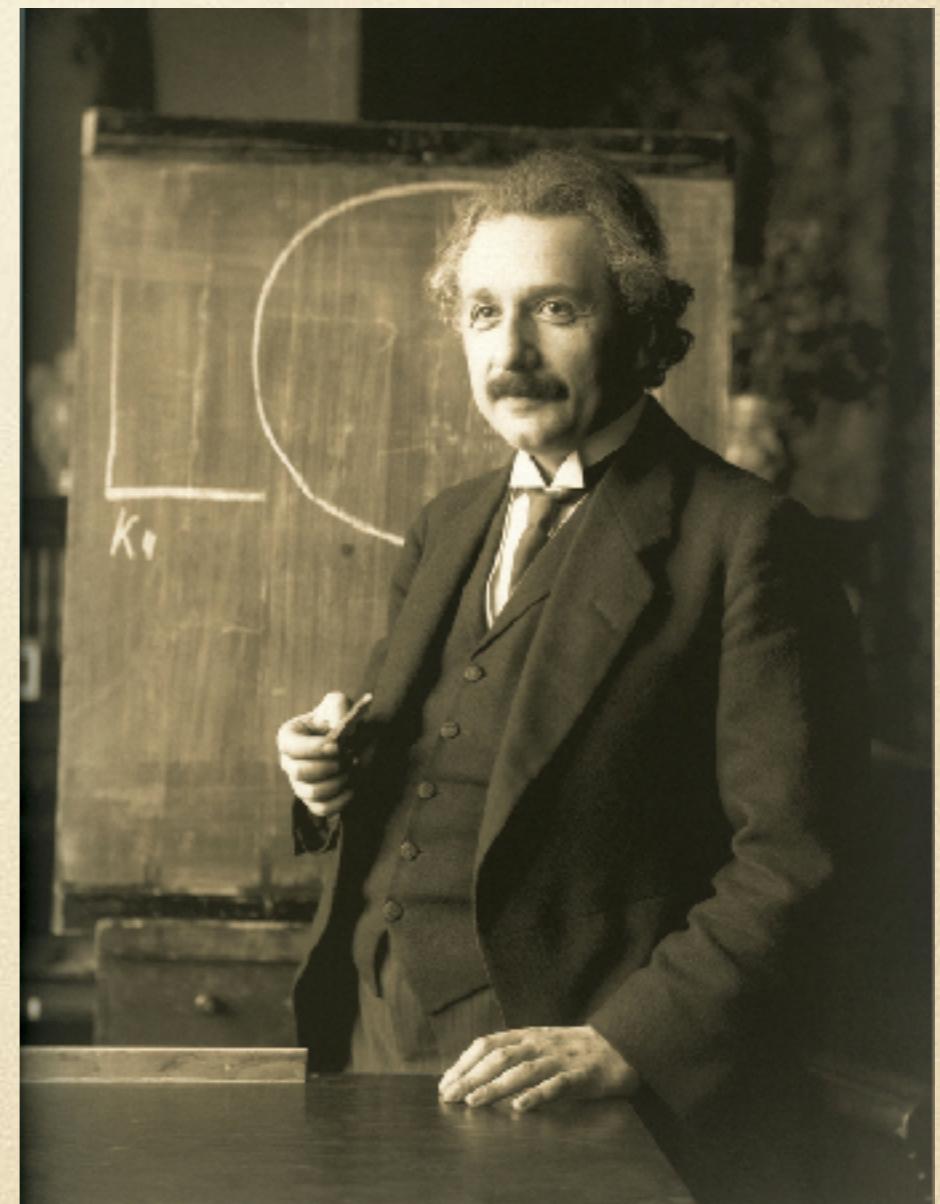


The CCP 1.0

The CCP 1.0

Albert Einstein

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$



NP 1921

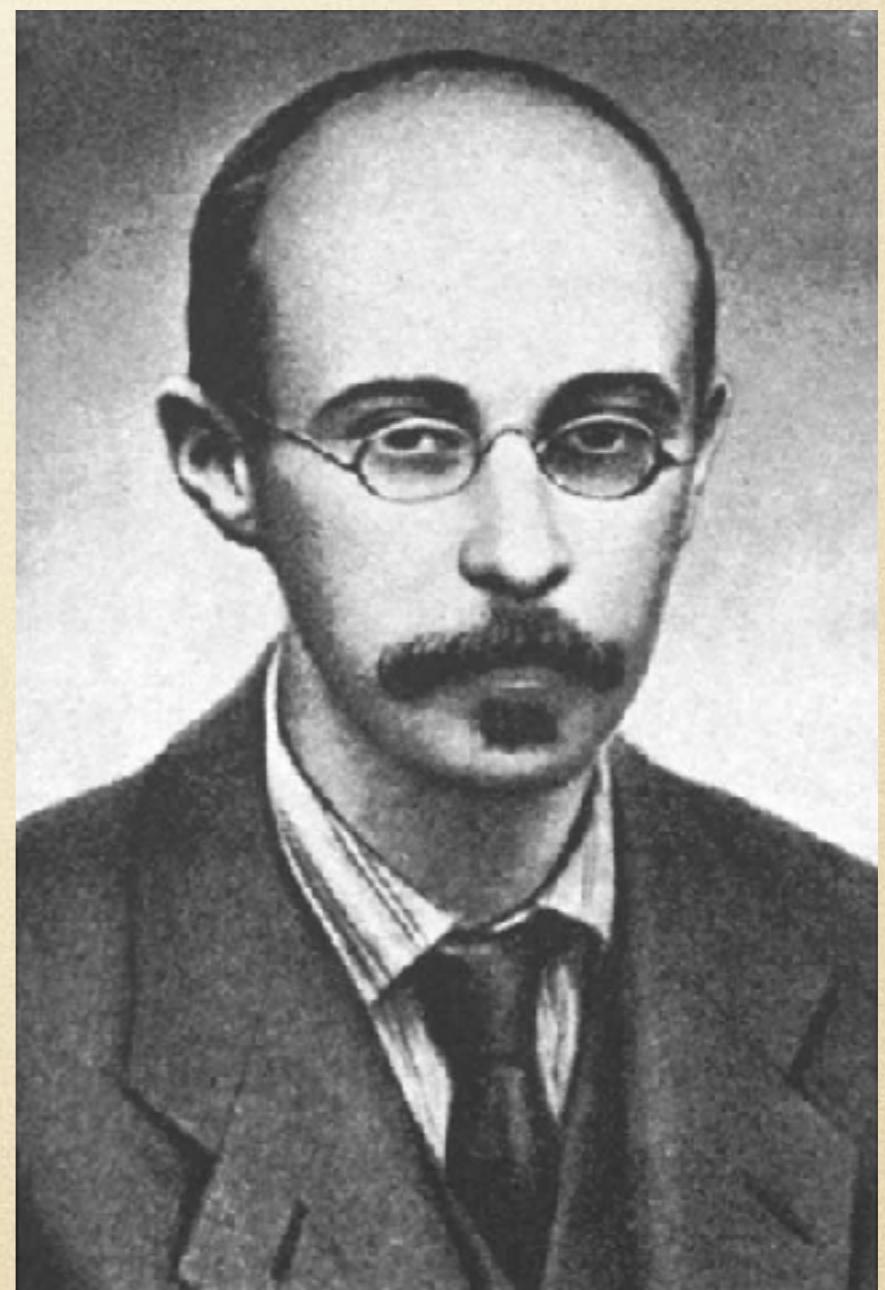
The CCP 1.0

Alexander Friedmann

$$ds^2 = a(t)ds_3^2 - dt^2$$

$$\frac{\dot{a}^2 + k}{a^3} = \frac{1}{3}8\pi G\rho$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p)$$



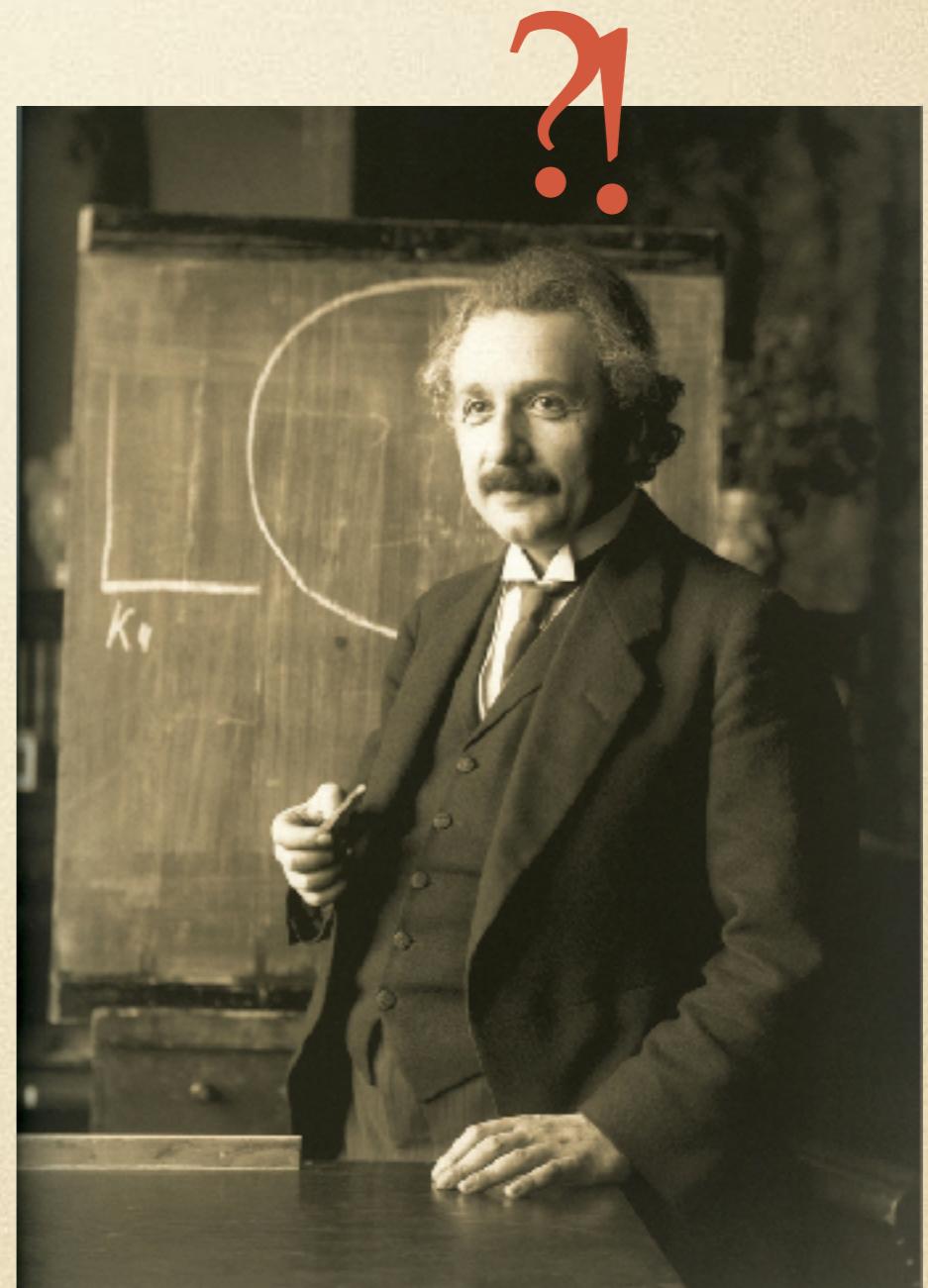
The CCP 1.0

Albert Einstein

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) \neq 0$$

$\Rightarrow \dot{a} \neq 0$ not static



The CCP 1.0

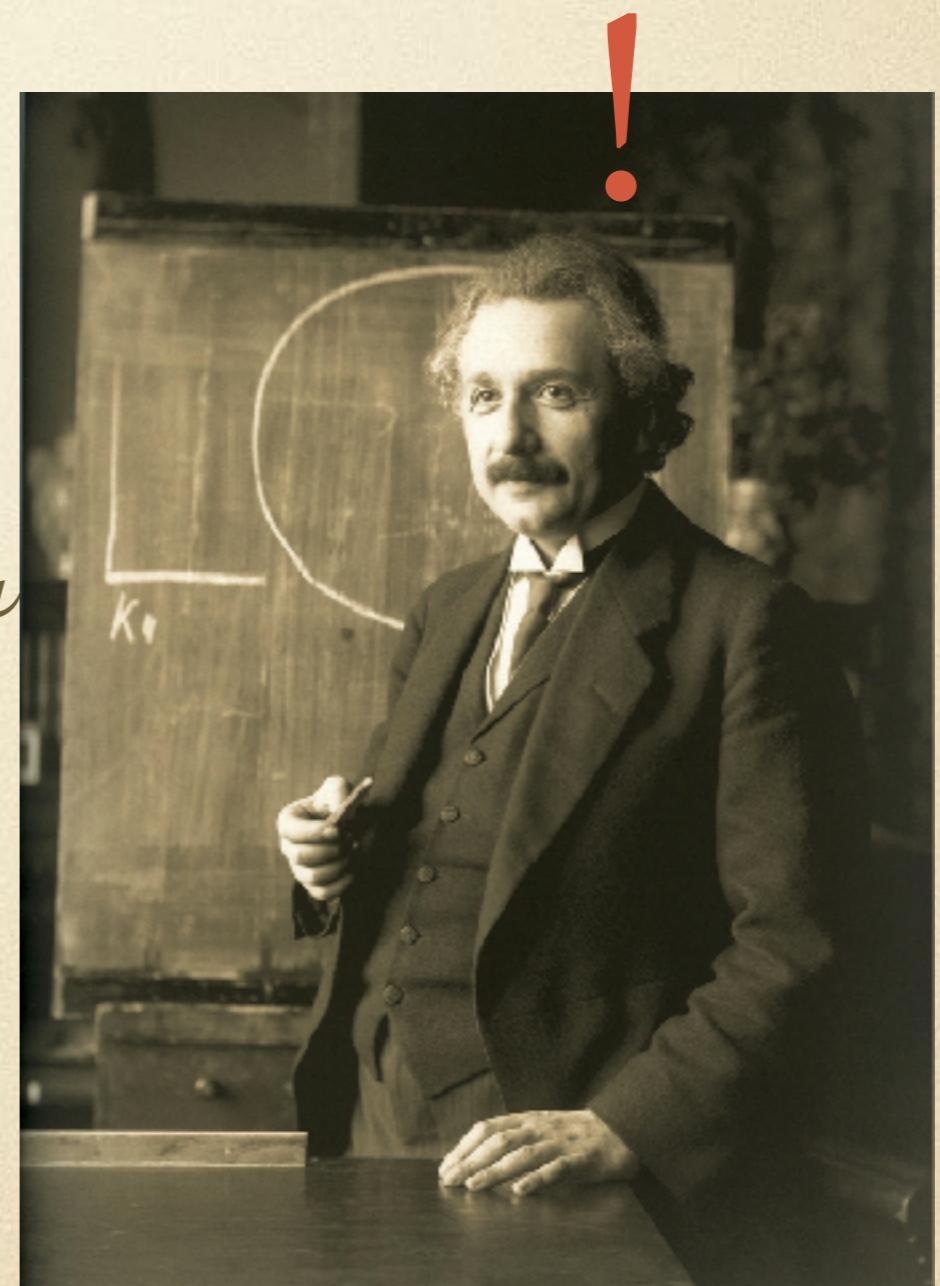
Albert Einstein

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \frac{\Lambda}{3} \equiv 0$$

!

static possible $\dot{a} \equiv 0$

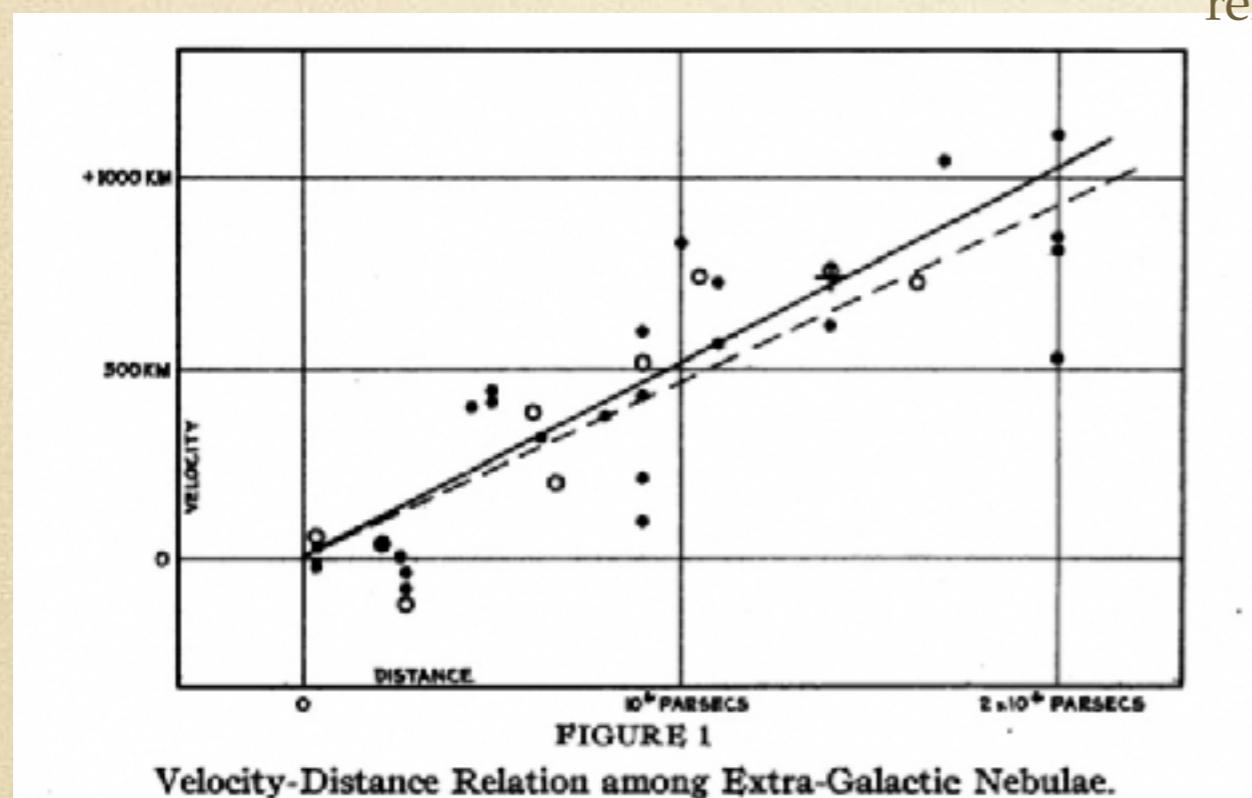


The CCP 2.0

The CCP 2.0

Edwin Hubble, Georges Lemaître
measurement:

ref [1]



1927,1929

The CCP 2.0

Edwin Hubble

measurement:

$$\dot{a} > 0 \quad \text{not static}$$

later:

$$\dot{a} = 67.66 \pm 0.42 \frac{\text{km/s}}{\text{Mpc}}$$

(Planck collaboration 2018)



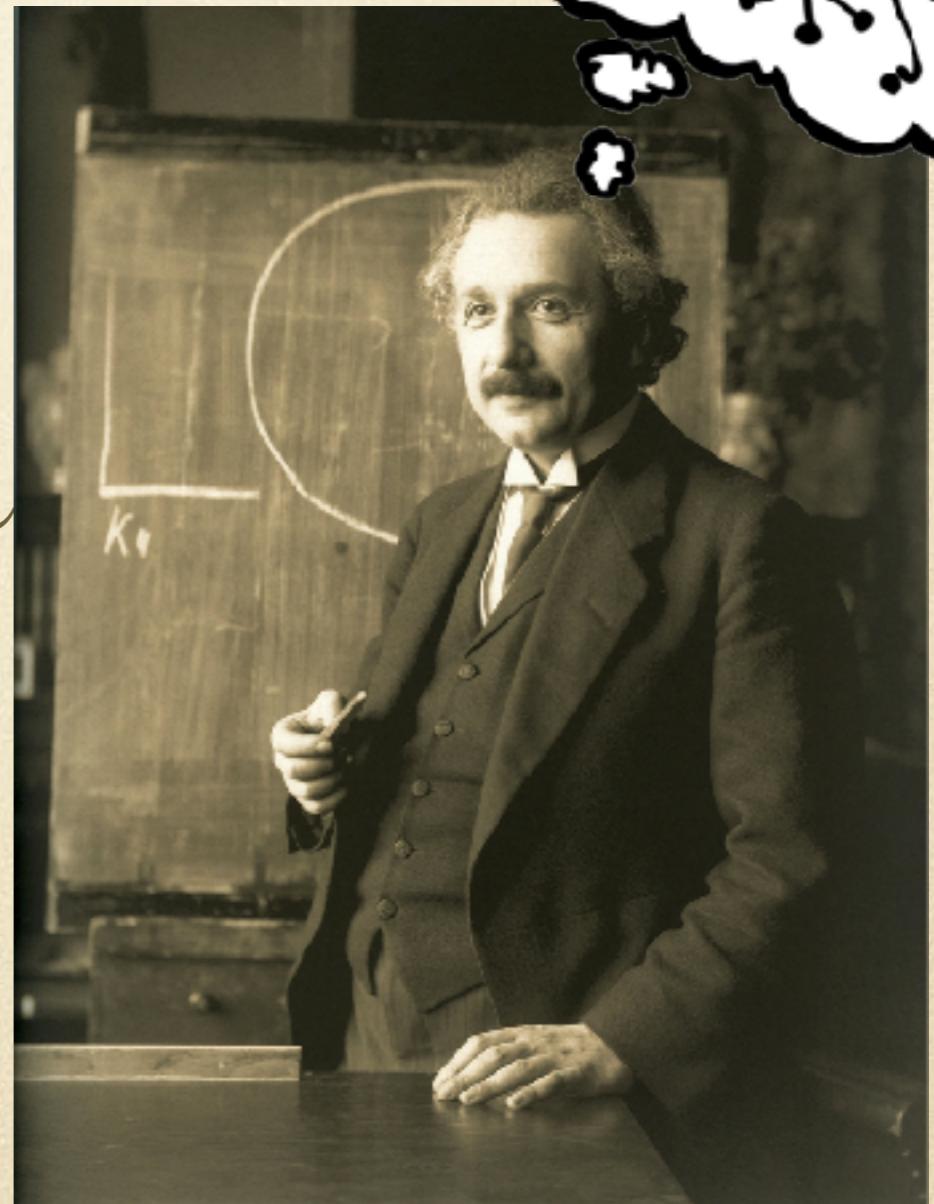
The CCP 2.0

Albert Einstein

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \frac{\Lambda}{a^2} = 0$$

not static $\dot{a} \neq 0$



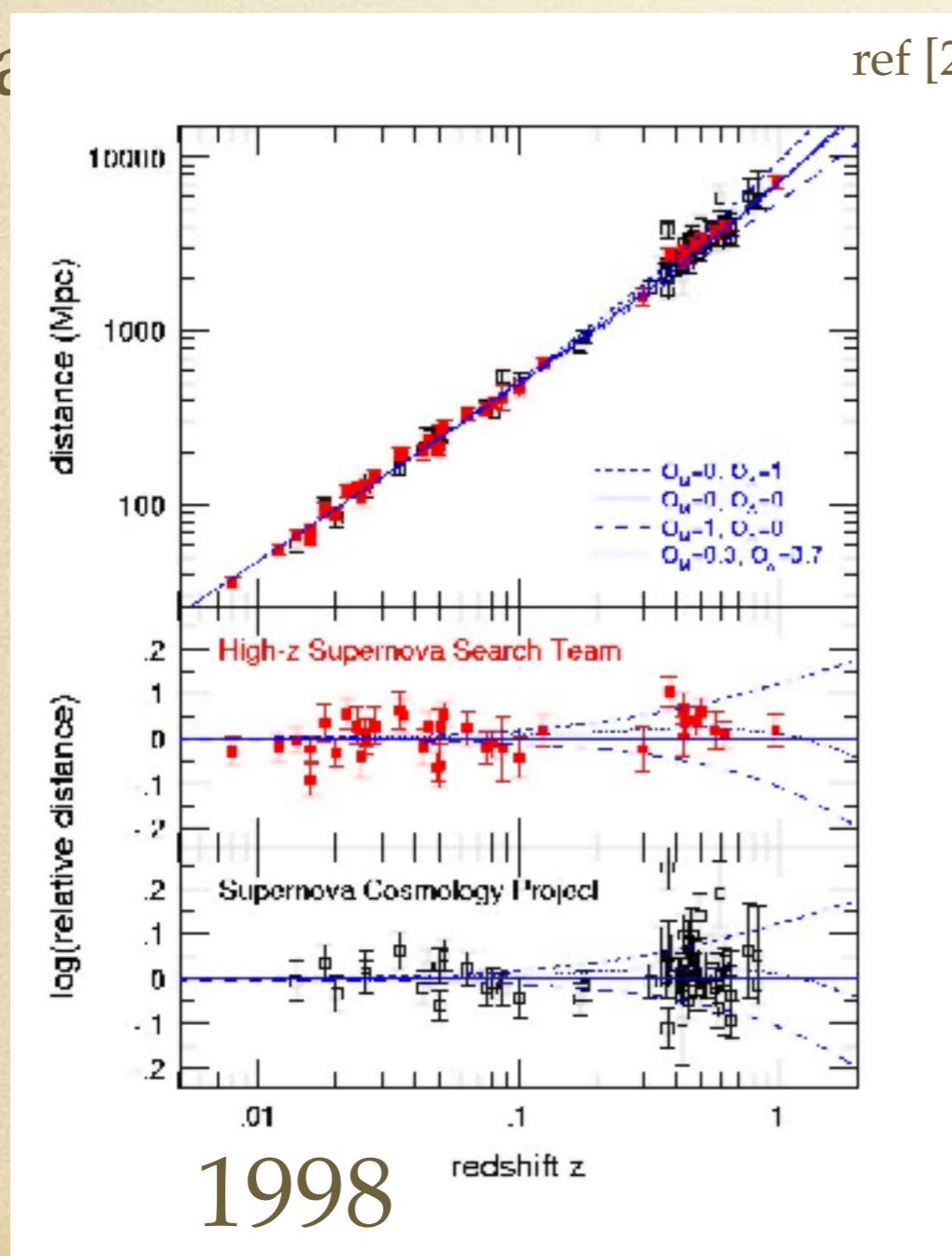
“biggest blunder”

The CCP 3.0

The CCP 3.0

S. Perlmutter, A. Riess, B. Schmidt, & others

mea



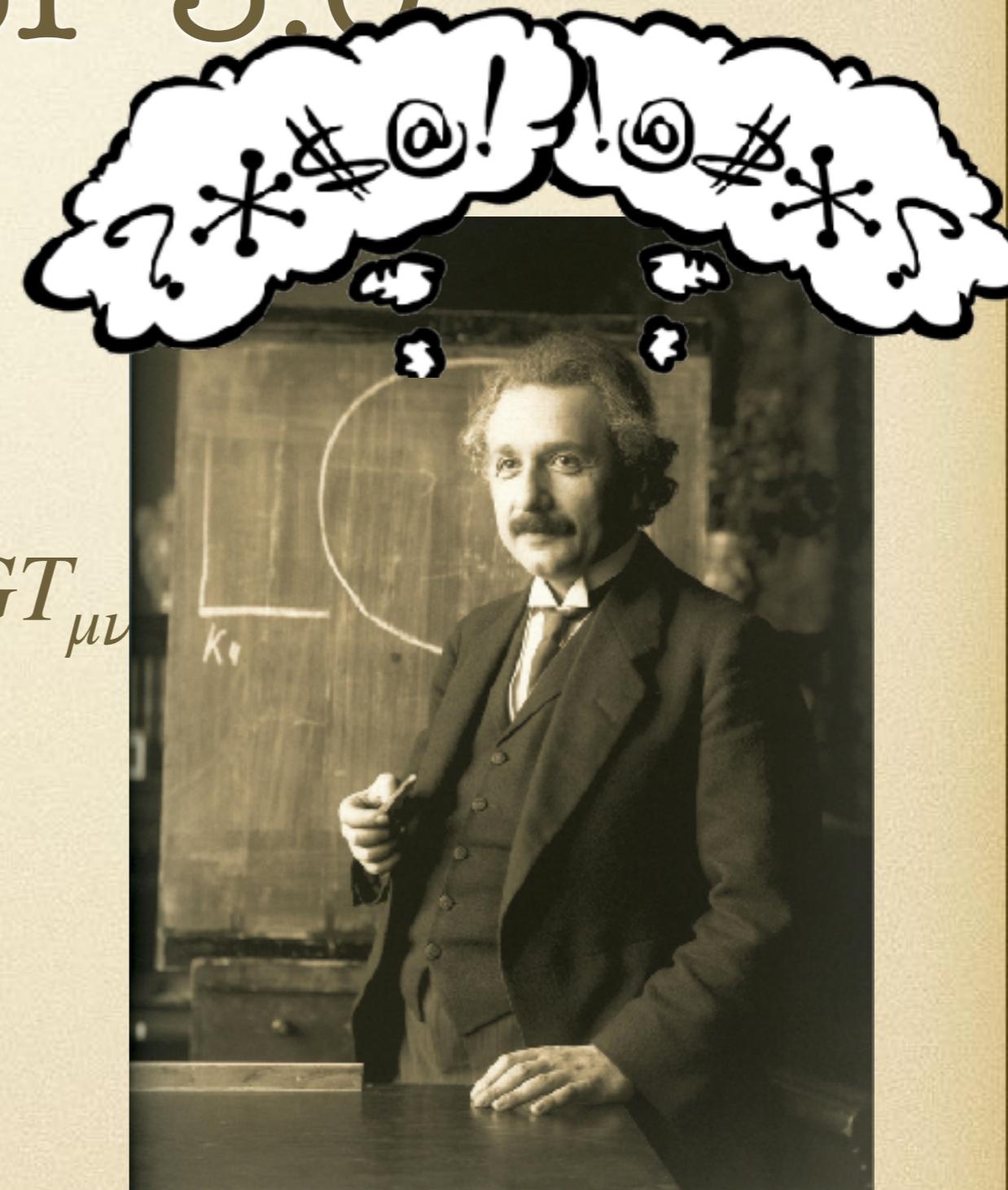
The CCP 3.0

Albert Einstein

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \Lambda$$

$$\Lambda > 0 \Rightarrow \ddot{a} > 0$$



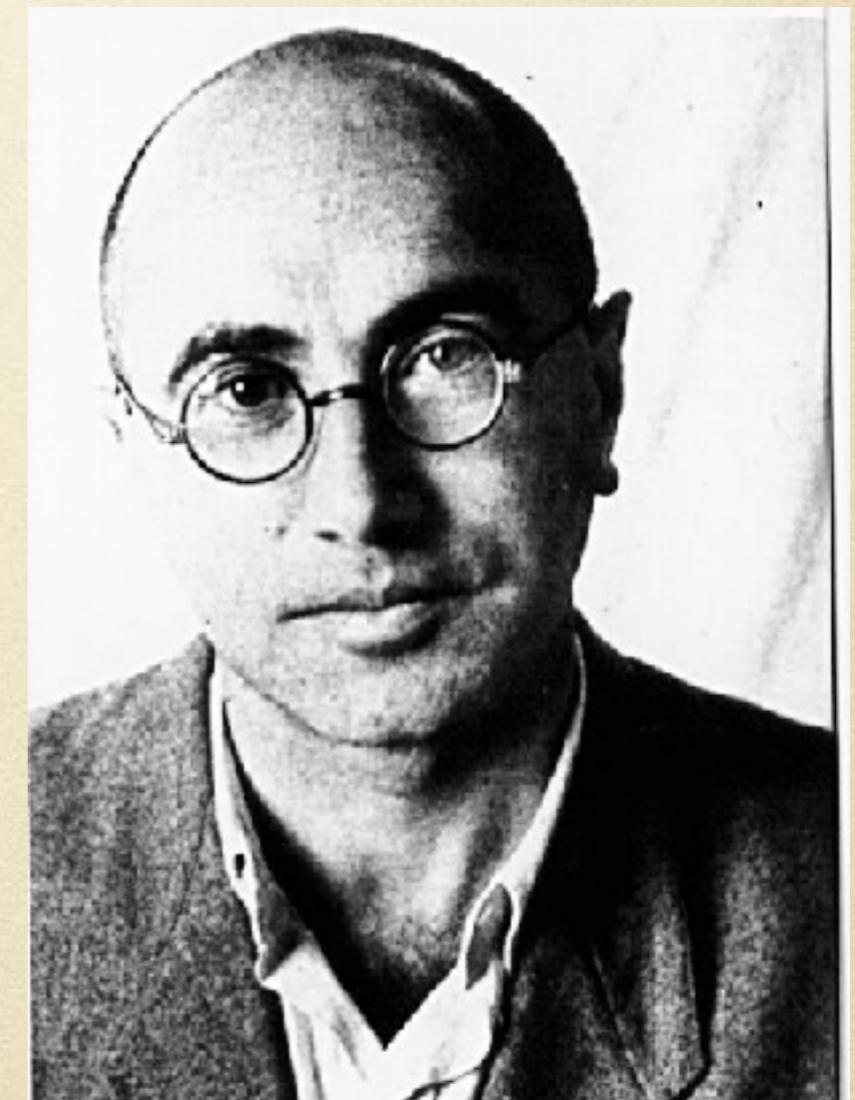
“???”

The CCP 3.0

Yakov Zeldovich

Quantum fluctuations
predict value of Λ

1967



The CCP 3.0

Steven Weinberg

Quantum fluctuations
predict value of Λ

Problem since 1998^{ref [3]}



The CCP 3.0

Quantum fluctuations
predict value of Λ

$$\left(\frac{\Lambda}{G}\right)_{eff} \sim M^4(1 + \dots)$$

ref [3]

effective quantum value

use:

different methods

cut-off mass scale of QFT

$G \sim \frac{1}{M_P^2}$

55

$$G \sim \frac{1}{M_P^2}$$

The CCP 3.0

Quantum fluctuations
predict value of Λ

effective
quantum value

$$(\Lambda)_{eff} \sim \frac{M^4}{M_P^2} (1 + \dots)$$

different methods

cut-off mass scale of QFT

Planck mass

The diagram illustrates the formula for the effective quantum value of Λ . It shows a horizontal red arrow pointing right, labeled "effective quantum value". Above this arrow is a mathematical expression: $(\Lambda)_{eff} \sim \frac{M^4}{M_P^2} (1 + \dots)$. To the left of the expression is the text "different methods" with a red arrow pointing towards it. To the right is the text "cut-off mass scale of QFT" with a red arrow pointing towards it. Below the expression is the text "Planck mass" with a red arrow pointing towards it.

The CCP 3.0

Quantum fluctuations
predict value of Λ

$$(\Lambda)_{eff} \sim \frac{M^4}{M_P^2} (1 + \dots)$$

Highest physical
mass scale

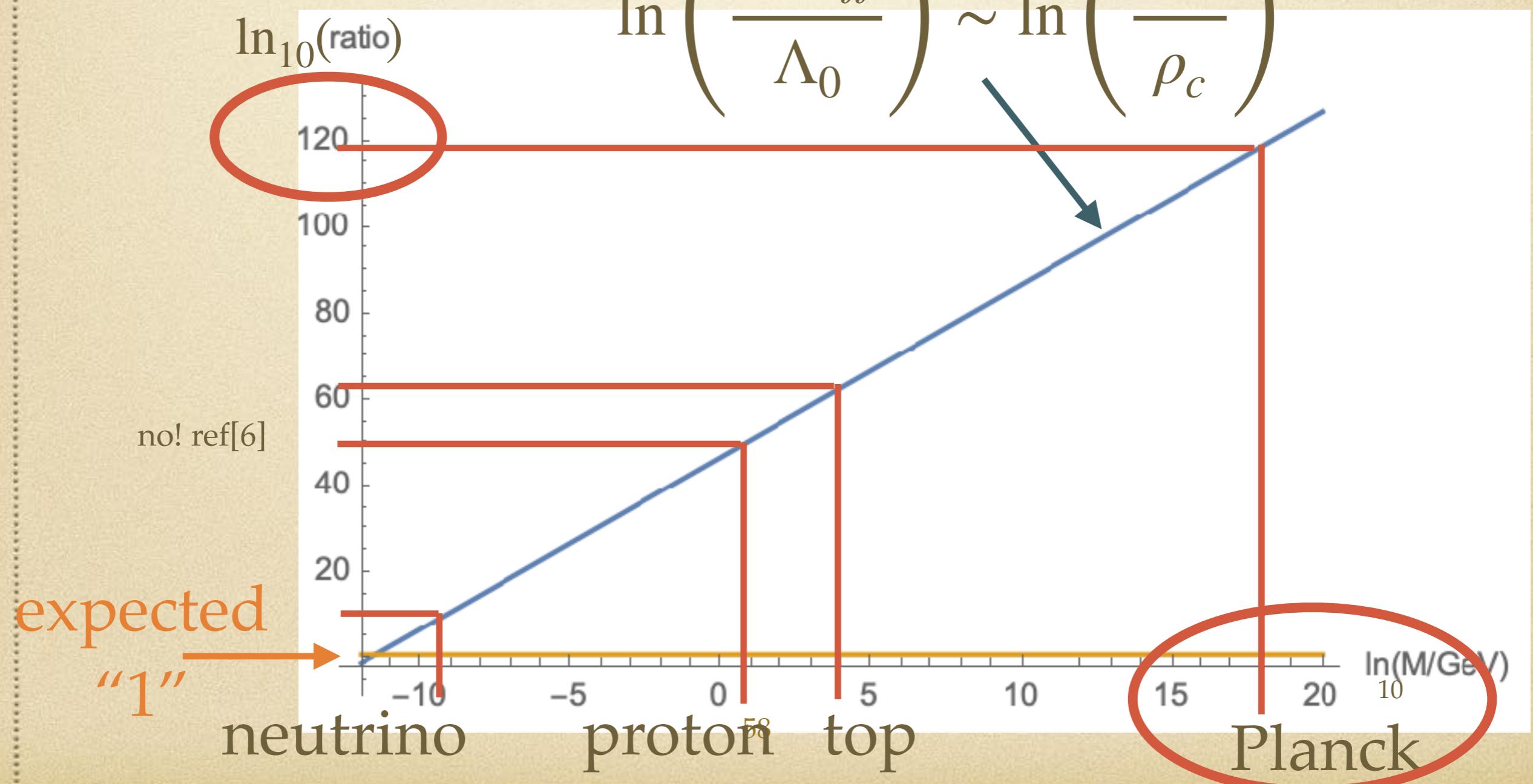
Observed value

$$\Lambda_o = \frac{\rho_c}{M_P^2} \approx \frac{10^{-47} GeV^4}{M_P^2}$$

observed critical
energy density

The CCP 3.0

Problem as a ratio:

$$\ln_{10}(\text{ratio}) \sim \ln\left(\frac{(\Lambda)_{eff}}{\Lambda_0}\right) \sim \ln\left(\frac{M^4}{\rho_c}\right)$$


The CCP 3.0

Problem as a ratio:

$$\frac{(\Lambda)_{eff}}{\Lambda_0} \sim \frac{1}{G_N \cdot \Lambda_0} \sim \frac{M_P^4}{\rho_c} \approx 10^{120}$$

we try to address
this problem

assuming there are quantum fluctuations
of gravity associated to the Planck scale

