

- $E^2 = \vec{p}^2 + m^2$
- $E = i\partial_t; p_j = -i\frac{\partial}{\partial x^j}$
- $-\frac{\partial^2}{\partial t^2}\varphi + \nabla^2\varphi - m^2\varphi = 0$, Klein-Gordon equation
- $\mathcal{L} = \frac{1}{2}\partial_\mu\varphi\partial^\mu\varphi - \frac{1}{2}m^2\varphi^2$
- Complex field: $\mathcal{L} = \partial_\mu\varphi^*\partial^\mu\varphi - \frac{1}{2}m^2\varphi^*\varphi$
- \mathcal{L} is invariant under $\varphi \rightarrow e^{i\alpha}\varphi, \dots$
- Conserved current (Noether theorem): $j^\mu = i(\varphi^*\partial^\mu\varphi - c.c.)$
- Probability density: $j^0 = i(\varphi^*\partial_t\varphi - c.c.)$. It can be negative!
- Solution (Pauli): j^μ is not a probability current but a charge current. Moreover φ must be a field, not a wave function.
- The K-G equation is the right relativistic equation to describe particles of spin=0.
- We use Sakita's signature: $(+, -, -, -)$

- To avoid negative probabilities, Dirac proposes that the equation must be first order in time derivative.
- Special relativity treats space and time equal, so the equation must be first order in spatial derivatives also.
- $H = \alpha \cdot p + \beta m$
- $H^2 = p \cdot p + m^2 = \frac{1}{2} \{ \alpha_i, \alpha_j \} p_i p_j + \{ \alpha_i, \beta \} p_i + \beta^2 m^2$
- $\frac{1}{2} \{ \alpha_i, \alpha_j \} = \delta_{ij}; \{ \alpha_i, \beta \} = 0; \beta^2 = 1$
- Dirac equation: $i \frac{\partial \psi}{\partial t} = -i \alpha \cdot \nabla \psi + \beta m \psi$
- Covariant form: $i \left(\beta \frac{\partial \psi}{\partial t} + \beta \alpha \cdot \nabla \psi \right) - m \psi = 0$
- $\gamma^0 = \beta; \gamma^i = \beta \alpha_i$
- $(i \gamma^\mu \partial_\mu - m) \psi = 0$
- $\bar{\psi} = \psi^\dagger \gamma^0; \mathcal{L} = \bar{\psi} (i \gamma^\mu \partial_\mu - m) \psi$
- Invariance: $\psi \rightarrow e^{i \alpha} \psi$
- Noether current: $j^\mu = \bar{\psi} \gamma^\mu \psi$: Probability density = $\psi^\dagger \psi \geq 0$