The Maxwell and Yang-Mills Gauge Fields

Gauge invariances

- a) Relativistic covariance is maintained
- b) The field equations do not determine certain longitudinal components
- c) We have constraints, that restrict the initial data

The classical degrees of freedom are the independent functions required as Initial data for the Cauchy problem of hyperbolic equations.

An elliptic equation $\nabla^2 \phi = 0$ Does not contain degrees of freedom

• Principle of minimal coupling

For a complex Dirac spinor of charge q

$$\Psi(x) \to \Psi'(x) \equiv e^{iq\theta(x)}\Psi(x)$$

Gauge field

$$A_{\mu}(x) \to A'_{\mu}(x) \equiv A_{\mu}(x) + \partial_{\mu}\theta(x)$$

Covariant derivative

$$D_{\mu}\Psi(x) \equiv (\partial_{\mu} - iqA_{\mu}(x))\Psi(x)$$

• Free gauge field

$$[D_{\mu}, D_{\nu}]\Psi = -ieqF_{\mu\nu}\Psi$$
$$F_{\mu\nu}(x) = \partial_{\mu}A_{\nu}(x) - \partial_{\nu}A_{\mu}(x)$$
$$E_{i} = F_{0i} \qquad \qquad B_{i} = \frac{1}{2}\varepsilon_{ijk}F_{jk}$$

Free equation of motion

 $\partial^{\mu}F_{\mu\nu} = 0$

Noether identity $\partial^{\nu}\partial^{\mu}F_{\mu\nu}=0$ signal of the gauge symmetry

$$egin{array}{lll} A_\mu o A_\mu + \partial_\mu heta \ Bianchi identity & \partial_\mu F_{
u
ho} + \partial_
u F_{
ho\mu} + \partial_
u F_{
ho\mu} = 0 \end{array}$$

Maxwell algebra

• For constant electromagnetic field there is a generalization of the Poincare group

$$[P_{\mu}, P_{\nu}] = i Z_{\mu\nu}$$

together with the generators of Lorentz transformations

Degrees of freedom

Gauge fixing, eg Coulomb gauge

 $\partial_i A_i(\vec{x}, t) = 0$

This condition does eliminate the gauge freedom

 $\begin{array}{ll} \partial_i A_i' = \partial^i (A_i + \partial_i \theta) = 0 & \quad \text{If} \quad \nabla^2 \theta = 0 \\ \text{which implies} \quad \theta(x) \equiv 0 \end{array}$

Maxwell equation

$$\nabla^2 A_0 - \partial_0 (\partial_i A_i) = 0$$

$$\Box A_i - \partial_i \partial_0 A_0 - \partial_i (\partial_j A_j) = 0$$

Degrees of freedom

In the Coulomb gauge

 $\nabla^2 A_0 = 0$ implies no degrees of freedom

 $\Box A_i = 0 \qquad \text{massless particle}$

Initial data $A_i(\vec{x},0), \ \dot{A}_i(\vec{x},0)$

2(D-2) independent degrees of freedom Off-shell degrees of freedom

(D-2) on-shell degrees of freedom helicity states

• The field strength verifies

 $\Box F_{\mu\nu} = 0$

Gauge invriant description that electromagnetic field describes massless particles

• Hamiltonian counting of degrees of freedom

Primary constraints $\pi^0 = 0$

Secondary constraint $\nabla \vec{\pi} = 0$

First class constraints $\pi^0 = 0$, $\nabla \vec{\pi} = 0$

Gauge fixing constraints $A^0 = 0$ $\nabla \vec{A} = 0$

Dirac bracket and number of degrees of freedom, 8-4=4=2+2

Noether identities

Combination of equation of motion that vanish identically

Action $S = S[\phi]$ equations of motion $\frac{\delta S}{\delta \phi^i}$ =0Gauge transformations $\delta \phi = R^i_{\alpha} \epsilon^{\alpha}$ δS δS

0.00

Variation of the action

$$\delta S\left[\phi\right] = \frac{\delta S}{\delta \phi^i} R^i_{\alpha} \epsilon^{\alpha}$$

Noether identities

$$\frac{\delta S}{\delta \phi^i} R^i_\alpha = 0$$

• QED

$$S[A_{\mu},\bar{\Psi},\Psi] = \int \mathrm{d}^{D}x \left[-\frac{1}{4}F^{\mu\nu}F_{\mu\nu} + \bar{\Psi}(\gamma^{\mu}D_{\mu} - m)\Psi\right]$$

$$\partial^{\mu}F_{\mu\nu} - ieq\bar{\Psi}\gamma_{\nu}\Psi = 0$$
.

Dual tensors

In D=4

$$\tilde{H}^{\mu\nu} \equiv -\frac{1}{2} \mathrm{i}\varepsilon^{\mu\nu\rho\sigma} H_{\rho\sigma}$$

Selfdual-antiselfdual

$$H^{\pm}_{\mu\nu} = \frac{1}{2} \left(H_{\mu\nu} \pm \tilde{H}_{\mu\nu} \right) , \qquad H^{\pm}_{\mu\nu} = \left(H^{\mp}_{\mu\nu} \right)^*$$
$$- \frac{1}{2} i \varepsilon_{\mu\nu}{}^{\rho\sigma} H^{\pm}_{\rho\sigma} = \pm H^{\pm}_{\mu\nu}$$

properties

$$G^{+\mu\nu}H^{-}{}_{\mu\nu} = 0, \qquad G^{\pm\rho(\mu}H^{\pm\nu)}{}_{\rho} = -\frac{1}{4}\eta^{\mu\nu}G^{\pm\rho\sigma}H^{\pm}{}_{\rho\sigma}, \qquad G^{+}{}_{\rho[\mu}H^{-}{}_{\nu]}{}^{\rho} = 0.$$

Duality for the electromagnetic field

• Free case

The Maxwell and Bianchi equations

 $\partial_{\mu}F^{\mu\nu} = 0, \qquad \partial_{\mu}\tilde{F}^{\mu\nu} = 0$

are invariant under the transformation

$$F^{\mu\nu} \to F'^{\mu\nu} = \mathrm{i}\tilde{F}^{\mu\nu}$$

property

$$F_{\mu\nu}\tilde{F}^{\mu\nu} = -\mathrm{i}\partial_{\mu}(\varepsilon^{\mu\nu\rho\sigma}A_{\nu}F_{\rho\sigma})$$

Chern-Simmons action L =A \wedge dA topological action

Equations of motion F=dA=0 flat connections