



Recent Developments in the Spectral Theory of Orthogonal Polynomials

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Lecture 1: Introduction and Overview

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- Lecture 3: Szegő-Widom asymptotics for Chebyshev Polynomials
- Lecture 4: Killip-Simon Theorems for Finite Gap Sets



References for Lecture 1

[OPUC] B. Simon, *Orthogonal Polynomials on the Unit Circle, Part 1: Classical Theory*, AMS Colloquium Series **54.1**, American Mathematical Society, Providence, RI, 2005.

[OPUC2] B. Simon, *Orthogonal Polynomials on the Unit Circle, Part 2: Spectral Theory*, AMS Colloquium Series, **54.2**, American Mathematical Society, Providence, RI, 2005.

[SzThm] B. Simon, *Szegő's Theorem and Its Descendants: Spectral Theory for L^2 Perturbations of Orthogonal Polynomials*, M. B. Porter Lectures, Princeton University Press, Princeton, NJ, 2011.

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Overview

This past year has seen three remarkable developments in the spectral theory of orthogonal and related polynomials.

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OPs

Orthogonal polynomials on the real line (OPRL) and on the unit circle (OPUC) are particularly useful because the inverse problems are easy—indeed the inverse problem is the OP definition as we'll see.

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OPs

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OPs also enter in many application—both specific polynomials and the general theory. Indeed, my own interest came from studying discrete Schrödinger operators on $\ell^2(\mathbb{Z})$

$$(Hu)_n = u_{n+1} + u_{n-1} + Vu_n$$

and the realization that when restricted to \mathbb{Z}_+ , one had a special case of OPRL.

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OPRL basics

μ will be a probability measure on \mathbb{R} . We'll always suppose that μ has bounded support $[a, b]$ which is not a finite set of points. (We then say that μ is non-trivial.) This implies that $1, x, x^2, \dots$ are independent since $\int |P(x)|^2 d\mu = 0 \Rightarrow \mu$ is supported on the zeroes of P .

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Apply Gram Schmidt to $1, x, \dots$ and get monic polynomials

$$P_j(x) = x^j + \alpha_{j,1}x^{j-1} + \dots$$

and orthonormal (ON) polynomials

$$p_j = P_j / \|P_j\|$$

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OPRL basics

More generally we can do the same for any probability measure of bounded support on \mathbb{C} .

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OPRL basics

More generally we can do the same for any probability measure of bounded support on \mathbb{C} .

One difference from the case of \mathbb{R} , the linear combination of $\{x^j\}_{j=0}^{\infty}$ are dense in $L^2(\mathbb{R}, d\mu)$ by Weierstrass. This may or may not be true if $\text{supp}(d\mu) \not\subset \mathbb{R}$.

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If $d\mu = d\theta/2\pi$ on $\partial\mathbb{D}$, the span of $\{z^j\}_{j=0}^{\infty}$ is not dense in L^2 (but is only H^2). Perhaps, surprisingly, as one can prove using Szegő's theorem, there are measures $d\mu$ on $\partial\mathbb{D}$ for which they are dense (e.g., μ purely singular).

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More significantly, the argument we'll give for our recursion relation fails if $\text{supp}(d\mu) \not\subset \mathbb{R}$.

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OPRL basics

Because $\langle P_j, xP_n \rangle = \langle xP_j, P_n \rangle$ (OPRL only)

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OPRL basics

Because $\langle P_j, xP_n \rangle = \langle xP_j, P_n \rangle$ (OPRL only) = 0 if $j < n - 1$, the P 's obey a three term recurrence relation:

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OPRL basics

Because $\langle P_j, xP_n \rangle = \langle xP_j, P_n \rangle$ (OPRL only) = 0 if $j < n - 1$, the P 's obey a three term recurrence relation: ($P_{-1} \equiv 0$); $\{a_j\}_{j=1}^{\infty}$, $\{b_j\}_{j=1}^{\infty}$: Jacobi recursion

$$xP_N = P_{N+1} + b_{N+1}P_N + a_N^2P_{N-1}$$

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$$xP_N = P_{N+1} + b_{N+1}P_N + a_N^2P_{N-1}$$

$$b_N \in \mathbb{R}, \quad a_N = \|P_N\|/\|P_{N-1}\|$$

These are called Jacobi parameters.

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These are called Jacobi parameters. This implies $\|P_N\| = a_N a_{N-1} \dots a_1$ (since $\|P_0\| = 1$).

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These are called Jacobi parameters. This implies $\|P_N\| = a_N a_{N-1} \dots a_1$ (since $\|P_0\| = 1$).

This, in turn, implies $p_n = P_n/a_1 \dots a_n$ obeys

$$xp_n = a_{n+1}p_{n+1} + b_{n+1}p_n + a_n p_{n-1}$$

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OPRL basics

We have thus solved the inverse problem, i.e., μ is the spectral data and $\{a_n, b_n\}_{n=1}^{\infty}$ are the descriptors of the object.

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In the orthonormal basis, $\{p_n\}_{n=0}^{\infty}$, multiplication by x has the matrix

$$J = \begin{pmatrix} b_1 & a_1 & 0 & 0 & \dots \\ a_1 & b_2 & a_2 & 0 & \dots \\ 0 & a_2 & b_3 & a_3 & \dots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

called a Jacobi matrix.

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Favard's Theorem

Since

$$b_n = \int x p_{n-1}^2(x) d\mu, \quad a_n = \int x p_{n-1}(x) p_n(x) d\mu$$

$$\text{supp}(\mu) \subset [-R, R] \Rightarrow |b_n| \leq R, |a_n| \leq R.$$

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Conversely, if $\sup_n (|a_n| + |b_n|) = \alpha < \infty$, J is a bounded matrix of norm at most 3α . In that case, the spectral theorem implies there is a measure $d\mu$ so that

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$$\langle (1, 0, \dots)^t, J^\ell (1, 0, \dots)^t \rangle = \int x^\ell d\mu(x)$$

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$$\langle (1, 0, \dots)^t, J^\ell (1, 0, \dots)^t \rangle = \int x^\ell d\mu(x)$$

If one uses Gram-Schmidt to orthonormalize $\{J^\ell (1, 0, \dots)^t\}_{\ell=0}^\infty$, one finds μ has Jacobi matrix exactly given by J .

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Favard's Theorem

We have thus proven Favard's Theorem (his paper was in 1935; really due to Stieltjes in 1894 or to Stone in 1932).

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Favard's Theorem. *There is a one–one correspondence between bounded Jacobi parameters*

$$\{a_n, b_n\}_{n=1}^{\infty} \in [(0, \infty) \times \mathbb{R}]^{\infty}$$

and non-trivial probability measures, μ , of bounded support via:

$$\mu \Rightarrow \{a_n, b_n\} \quad (\text{OP recursion})$$

$$\{a_n, b_n\} \Rightarrow \mu \quad (\text{Spectral Theorem})$$

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$$\{a_n, b_n\} \Rightarrow \mu \quad (\text{Spectral Theorem})$$

There are also results for μ 's with unbounded support so long as $\int x^n d\mu < \infty$. In this case, $\{a_n, b_n\} \Rightarrow \mu$ may not be unique because J may not be essentially self-adjoint on vectors of finite support.

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OPUC basics

Let $d\mu$ be a non-trivial probability measure on $\partial\mathbb{D}$. As in the OPRL case, we use Gram-Schmidt to define monic OPs, $\Phi_n(z)$ and ON OP's $\varphi_n(z)$.

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In the OPRL case, if z is multiplication by the underlying variable, $z^* = z$. This is replaced by $z^*z = 1$.

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In the OPRL case, if z is multiplication by the underlying variable, $z^* = z$. This is replaced by $z^*z = 1$.

In the OPRL case, $P_{n+1} - xP_n \perp \{1, x_1, \dots, x_{n-2}\}$.

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In the OPRL case, if z is multiplication by the underlying variable, $z^* = z$. This is replaced by $z^*z = 1$.

In the OPRL case, $P_{n+1} - xP_n \perp \{1, x_1, \dots, x_{n-2}\}$. In the OPUC case, $\Phi_{n+1} - z\Phi_n \perp \{z, \dots, z^n\}$, since

$$\langle z\Phi_n, z^j \rangle = \langle \Phi_n, z^{j-1} \rangle$$

if $j \geq 1$.

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In the OPRL case, if z is multiplication by the underlying variable, $z^* = z$. This is replaced by $z^*z = 1$.

In the OPRL case, $P_{n+1} - xP_n \perp \{1, x_1, \dots, x_{n-2}\}$. In the OPUC case, $\Phi_{n+1} - z\Phi_n \perp \{z, \dots, z^n\}$, since

$$\langle z\Phi_n, z^j \rangle = \langle \Phi_n, z^{j-1} \rangle$$

if $j \geq 1$.

In the OPRL case, we used $\deg P = n$ and $P \perp \{1, x, \dots, x^{n-2}\} \Rightarrow P = c_1P_n + c_2P_{n-1}$.

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Let $d\mu$ be a non-trivial probability measure on $\partial\mathbb{D}$. As in the OPRL case, we use Gram-Schmidt to define monic OPs, $\Phi_n(z)$ and ON OP's $\varphi_n(z)$.

In the OPRL case, if z is multiplication by the underlying variable, $z^* = z$. This is replaced by $z^*z = 1$.

In the OPRL case, $P_{n+1} - xP_n \perp \{1, x_1, \dots, x_{n-2}\}$. In the OPUC case, $\Phi_{n+1} - z\Phi_n \perp \{z, \dots, z^n\}$, since

$$\langle z\Phi_n, z^j \rangle = \langle \Phi_n, z^{j-1} \rangle$$

if $j \geq 1$.

In the OPRL case, we used $\deg P = n$ and $P \perp \{1, x, \dots, x^{n-2}\} \Rightarrow P = c_1P_n + c_2P_{n-1}$.

In the OPUC case, we want to characterize $\deg P = n$, $P \perp \{z, z^2, \dots, z^n\}$.



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Define $*$ on degree n polynomials to themselves by

$$Q^*(z) = z^n \overline{Q\left(\frac{1}{\bar{z}}\right)}$$

(bad but standard notation!) or

$$Q(z) = \sum_{j=0}^n c_j z^j \Rightarrow Q^*(z) = \sum_{j=0}^n \bar{c}_{n-j} z^j$$

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$$Q(z) = \sum_{j=0}^n c_j z^j \Rightarrow Q^*(z) = \sum_{j=0}^n \bar{c}_{n-j} z^j$$

Then, $*$ is anti-unitary and so for $\deg Q = n$

$$Q \perp \{1, \dots, z^{n-1}\} \Leftrightarrow Q = c \Phi_n$$

is equivalent to

$$Q \perp \{z, \dots, z^n\} \Leftrightarrow Q = c \Phi_n^*$$

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Szegő recursion and Verblunsky coefficients

Thus, we see, there are parameters $\{\alpha_n\}_{n=0}^{\infty}$ (called Verblunsky coefficients) so that

$$\Phi_{n+1}(z) = z\Phi_n - \bar{\alpha}_n\Phi_n^*(z)$$

This is the Szegő Recursion (History: Szegő and Geronimus in 1939; Verblunsky in 1935–36)

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Applying $*$ for deg $n + 1$ polynomials to this yields

$$\Phi_{n+1}^*(z) = \Phi_n^*(z) - \alpha_n z\Phi_n$$

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The strange looking $-\bar{\alpha}_n$ rather than say $+\alpha_n$ is to have the α_n be the Schur parameter of the Schur function of μ (Geronimus); also the Verblunsky parameterization then agrees with α_n . These are discussed in [OPUC1].

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For OPRL, we saw $\|P_{n+1}\|/\|P_n\| = a_{n+1}$. We are looking for the analog for OPUC.

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$$\text{Szegő Recursion} \Rightarrow \Phi_{n+1} + \bar{\alpha}_n \Phi_n^* = z\Phi_n$$

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Szegő Recursion $\Rightarrow \Phi_{n+1} + \bar{\alpha}_n \Phi_n^* = z\Phi_n$

$$\Phi_{n+1} \perp \Phi_n^* \Rightarrow \|\Phi_{n+1}\|^2 + |\alpha_n|^2 \|\Phi_n^*\|^2 = \|z\Phi_n\|^2$$

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Multiplication by z unitary plus $*$ antiunitary \Rightarrow

$$\|\Phi_{n+1}\|^2 = \rho_n^2 \|\Phi_n\|^2; \quad \rho_n^2 = 1 - |\alpha_n|^2$$

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$$\|\Phi_{n+1}\|^2 = \rho_n^2 \|\Phi_n\|^2; \quad \rho_n^2 = 1 - |\alpha_n|^2$$

which implies $|\alpha_n| < 1$ (i.e., $\alpha_n \in \mathbb{D}$) and

$$\|\Phi_n\| = \rho_{n-1} \cdots \rho_0$$

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Verblunsky's Theorem

Verblunsky's Theorem. *There is a one–one correspondence between Verblunsky coefficients*

$$\{\alpha_n\}_{n=0}^{\infty} \in \mathbb{D}^{\infty}$$

and non-trivial probability measures, μ , of supported on \mathbb{D} .

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The measure determines the α 's via forming OPUC and looking at recursion relations.

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The measure determines the α 's via forming OPUC and looking at recursion relations. One way of going in the opposite direction is to write multiplication by z in an ON basis obtained by orthonormalizing $\{1, z, z^{-1}, z^2, z^{-2}, \dots\}$. This basis can be written in terms of φ_n and φ_n^* .

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For Jacobi, it is similar.

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For Jacobi, it is similar. Instead of $\rho_{n-1} = 0$, we have $a_n = 0$ and there are again $2n - 1$ free parameters $\{a_j\}_{j=1}^{n-1} \cup \{b_j\}_{j=1}^n$.

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Variational Form

We turn next to Szegő's Theorem. In 1914, while still a student in Budapest, Szegő proved a conjecture of Polya on the asymptotics of determinants of Toeplitz matrices.

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By a simple argument, $\|\Phi_n\|^2$ is the minimum of $\int |P(x)|^2 d\mu(x)$ over all monic polynomials of degree n . Thus $\|\Phi_n\|^2 \leq \|z\Phi_{n-1}\|^2 = \|\Phi_{n-1}\|^2$ so the norms are decreasing and thus have a limit. Szegő identified this limit:

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Szegő's Theorem

$$d\mu = w(\theta) \frac{d\theta}{2\pi} + d\mu_s \Rightarrow \lim_{n \rightarrow \infty} \|\Phi_n\|^2 = \exp \left(\int \log(w(\theta)) \frac{d\theta}{2\pi} \right)$$

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Szegő's Theorem as a Sum Rule

The version we just stated is from Szegő's 1920-21 paper on the foundations of OPUC which included the connection of OPUC to Toeplitz matrices.

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The version we just stated is from Szegő's 1920-21 paper on the foundations of OPUC which included the connection of OPUC to Toeplitz matrices. He also only included the case $\mu_s = 0$, in part because at that time singular measures were not very well known or studied.

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$$\sum_{n=0}^{\infty} \log(1 - |\alpha_n|^2) = \int \log(w(\theta)) \frac{d\theta}{2\pi}$$

This has been called the *Szegő-Verblunsky sum rule* (by Gamboa et. al in the work I'll discuss in Lecture 2).

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Szegő's Theorem as a Sum Rule

The version we just stated is from Szegő's 1920-21 paper on the foundations of OPUC which included the connection of OPUC to Toeplitz matrices. He also only included the case $\mu_s = 0$, in part because at that time singular measures were not very well known or studied. Szegő only found the recursion relation in 1939, so it didn't have the form that relies on Verblunsky coefficients.

Verblunsky, in 1935, found his parameters (with a different definition), allowed $\mu_s \neq 0$ and wrote in Szegő's theorem in the form:

$$\sum_{n=0}^{\infty} \log(1 - |\alpha_n|^2) = \int \log(w(\theta)) \frac{d\theta}{2\pi}$$

This has been called the *Szegő-Verblunsky sum rule* (by Gamboa et. al in the work I'll discuss in Lecture 2). It is a precursor of the KdV sum rules. Verblunsky's work was widely ignored until about 15 years ago.

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Szegő Theorem Gem

Much of spectral theory involves studying the relations between spectral information and parameters of the equations defining the spectral information.

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Szegő Theorem Gem

Much of spectral theory involves studying the relations between spectral information and parameters of the equations defining the spectral information. The most common is that information on the parameters of the equation implies information about the spectrum,

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Szegő Theorem Gem

Much of spectral theory involves studying the relations between spectral information and parameters of the equations defining the spectral information. The most common is that information on the parameters of the equation implies information about the spectrum, e.g. Lieb-Thirring bounds that say L^p properties of the potential yield information on the negative eigenvalues.

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Less common is information in the other direction.

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Less common is information in the other direction. I invented the name “gems of spectral theory” for equivalences.

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Sum rules where one side only involves spectral information and one side parameters generate gems by noting the equivalence of the two sides being finite.

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Sum rules where one side only involves spectral information and one side parameters generate gems by noting the equivalence of the two sides being finite. For this to work easily, both sides have to be positive (or negative) to prevent the issue of cancelling infinities.

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$$\sum_{n=0}^{\infty} |\alpha_n|^2 \iff \int \log(w(\theta)) \frac{d\theta}{2\pi} > -\infty$$

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$$\sum_{n=0}^{\infty} |\alpha_n|^2 \iff \int \log(w(\theta)) \frac{d\theta}{2\pi} > -\infty$$

This is especially interesting because it implies there exists ℓ^2 α 's with essentially arbitrary imbedded singular spectrum!

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Szegő's Calculation

Next, I want to discuss the issue of Szegő Asymptotics.

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Szegő's Calculation

Next, I want to discuss the issue of Szegő Asymptotics. If $\int \log(w(\theta)) \frac{d\theta}{2\pi} > -\infty$, one can define a function, $D(z)$ on \mathbb{D} by

$$D(z) = \exp\left(\int \frac{e^{i\theta} + z}{e^{i\theta} - z} \log(w(\theta)) \frac{d\theta}{4\pi}\right)$$

By a cutoff argument, $D \in H^2(\mathbb{D})$ and $|D(e^{i\theta})|^2 = f(\theta)$ in terms of boundary values.

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By a cutoff argument, $D \in H^2(\mathbb{D})$ and $|D(e^{i\theta})|^2 = f(\theta)$ in terms of boundary values. We have the following beautiful calculation of Szegő:

$$\int |\varphi_n^*(e^{i\theta}) D(e^{i\theta}) - 1|^2 \frac{d\theta}{2\pi} + \int |\varphi_n^*(e^{i\theta})|^2 d\mu_s = 2(1 - \prod_{j=n}^{\infty} \rho_j)$$

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$$\text{LHS} = \int \frac{d\theta}{2\pi} + \int |\varphi_n^*(e^{i\theta})|^2 d\mu - 2 \operatorname{Re} \int D(e^{i\theta}) \varphi_n^*(e^{i\theta}) \frac{d\theta}{2\pi}$$

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$$\begin{aligned} \text{LHS} &= \int \frac{d\theta}{2\pi} + \int |\varphi_n^*(e^{i\theta})|^2 d\mu - 2 \operatorname{Re} \int D(e^{i\theta}) \varphi_n^*(e^{i\theta}) \frac{d\theta}{2\pi} \\ &= 2 - 2 \operatorname{Re}(D(0) \varphi_n^*(0)) \end{aligned}$$

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$$= 2 \left[1 - \prod_{j=0}^{\infty} \rho_j \left(\prod_{j=0}^{n-1} \rho_j^{-1} \right) \right]$$

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Since $\text{RHS} \rightarrow 0$ as $n \rightarrow \infty$ (if the product converges, i.e., if the Szegő condition holds), each term goes to zero.

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Thus $\int |\varphi_n^*(e^{i\theta})|^2 d\mu_s \rightarrow 0$ and $\varphi_n^* D \rightarrow 1$ in $L^2(\partial\mathbb{D}, \frac{d\theta}{2\pi})$.

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Since the Poisson kernel $P_z(e^{i\theta})$ is L^2 uniformly for $|z| \leq r < 1$, $\varphi_n^*(z) D(z) \rightarrow 1$ uniformly on $|z| \leq r < 1$.

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Thus, uniformly in $|z| \geq r^{-1} > 1$,

$$z^{-n} \varphi_n(z) \rightarrow \left[\overline{D\left(\frac{1}{\bar{z}}\right)} \right]^{-1}$$

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Thus, uniformly in $|z| \geq r^{-1} > 1$,

$$z^{-n} \varphi_n(z) \rightarrow \left[\overline{D\left(\frac{1}{\bar{z}}\right)} \right]^{-1}$$

which is called *Szegő asymptotics* for φ_n .

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Szegő's Asymptotics for $[-2, 2]$

In 1922, Szegő found a way to write OPRL for measures on $[-2, 2]$ in terms of OPUC for the measure dragged to $\partial\mathbb{D}$ using the map $e^{i\theta} \mapsto x = 2 \cos(\theta)$.

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$$\int_{-2}^2 (4 - x^2)^{-1/2} \log(f(x)) dx > -\infty$$

called the Szegő condition for $[-2, 2]$.

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called the Szegő condition for $[-2, 2]$. The weight $(4 - x^2)^{-1/2}$ comes from $d\theta = (4 - x^2)^{-1/2} dx$.

Szegő Asymptotics translates to the existence of analytic functions, G , on \mathbb{D} and, \tilde{G} , on $\mathbb{C} \setminus [-2, 2]$ so that if the Szegő condition for $[-2, 2]$ holds.

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Szegő's Asymptotics for $[-2, 2]$

$$z^{-n} P_n \left(z + \frac{1}{z} \right) \rightarrow G(z)$$

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Szegő's Asymptotics for $[-2, 2]$

$$z^{-n} P_n \left(z + \frac{1}{z} \right) \rightarrow G(z)$$

Equivalently, for $x \in \mathbb{C} \setminus [-2, 2]$

$$\left(\frac{x}{2} + \sqrt{\left(\frac{x}{2} \right)^2 - 1} \right)^{-n} P_n(x) \rightarrow \tilde{G}(x)$$

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Various authors allowed adding discrete point spectrum outside $[-2, 2]$ until around 2000, Peherdsdorfer-Yuditskii and Killip-Simon got the ultimate result in cases where the Szegő condition for $[-2, 2]$ holds.

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$$\left(\frac{x}{2} + \sqrt{\left(\frac{x}{2}\right)^2 - 1} \right)^{-n} P_n(x) \rightarrow \tilde{G}(x)$$

Various authors allowed adding discrete point spectrum outside $[-2, 2]$ until around 2000, Peherdsdorfer-Yuditskii and Killip-Simon got the ultimate result in cases where the Szegő condition for $[-2, 2]$ holds. In 2006, Damanik-Simon found necessary and sufficient condition on the Jacobi parameters for Szegő Asymptotics to hold that, surprisingly to some, included some for which the Szegő condition for $[-2, 2]$ fails.

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Szegő's Asymptotics for $[-2, 2]$

The function $x \mapsto \frac{x}{2} + \sqrt{\left(\frac{x}{2}\right)^2 - 1}$ that replaces z when one moves Szegő Asymptotics from \mathbb{D} to $[-2, 2]$ can be understood by noting that its \log is the potential theorist's Green's function,

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- 1 It is harmonic and positive on $\mathbb{C} \setminus e$

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- 1 It is harmonic and positive on $\mathbb{C} \setminus \epsilon$
- 2 For q.e. $x \in \epsilon$, the boundary value is 0

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The function $x \mapsto \frac{x}{2} + \sqrt{\left(\frac{x}{2}\right)^2 - 1}$ that replaces z when one moves Szegő Asymptotics from \mathbb{D} to $[-2, 2]$ can be understood by noting that its \log is the potential theorist's Green's function, that is the unique function on \mathbb{C} (with $\epsilon = [-2, 2]$)

- 1** It is harmonic and positive on $\mathbb{C} \setminus \epsilon$
- 2** For q.e. $x \in \epsilon$, the boundary value is 0
- 3** $G_\epsilon(z) - \log |z|$ is harmonic at ∞

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Szegő's Asymptotics for $[-2, 2]$

The function $x \mapsto \frac{x}{2} + \sqrt{\left(\frac{x}{2}\right)^2 - 1}$ that replaces z when one moves Szegő Asymptotics from \mathbb{D} to $[-2, 2]$ can be understood by noting that its \log is the potential theorist's Green's function, that is the unique function on \mathbb{C} (with $\epsilon = [-2, 2]$)

- 1 It is harmonic and positive on $\mathbb{C} \setminus \epsilon$
- 2 For q.e. $x \in \epsilon$, the boundary value is 0
- 3 $G_\epsilon(z) - \log |z|$ is harmonic at ∞

Moreover, near ∞

$$\exp(G_\epsilon(z)) = \frac{|z|}{C(\epsilon)} + O(1)$$

where $C(\epsilon)$ is the logarithmic capacity of ϵ which is 1 for $\epsilon = [-2, 2]$.

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The Gem

In 2000, Rowan Killip and I proved the following OPRL analog of Szegő's Theorem.

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In 2000, Rowan Killip and I proved the following OPRL analog of Szegő's Theorem.

Killip–Simon Theorem *Let $d\mu(x) = f(x) dx + d\mu_s$ with Jacobi parameters $\{a_n, b_n\}_{n=1}^{\infty}$. Then*

$$\sum_{n=1}^{\infty} (a_n - 1)^2 + b_n^2 < \infty$$

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$$\sum_{n=1}^{\infty} (a_n - 1)^2 + b_n^2 < \infty$$

if and only if

(i) (Blumental–Weyl) $\sigma_{\text{ess}}(J) = \text{ess supp}(d\mu) = [-2, 2]$, i.e., $\text{supp}(d\mu)$ is a set of pure points whose only possible limit points are ± 2 : $E_1^- < E_2^- < \dots < -2$; $2 < \dots < E_2^+ < E_1^+$.

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(ii) (Lieb–Thirring) $\sum_{\pm, j} (|E_j^\pm| - 2)^{3/2} < \infty$.

(iii) (Quasi-Szegő) $\int (x^2 - 4)^{1/2} \log(f(x)) dx < \infty$.

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The Gem

If J_0 is Jacobi matrix, $a_n \equiv 1$, $b_n \equiv 0$, the L^2 condition is

$$\mathrm{Tr}((J - J_0)^2) < \infty$$

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For Schrödinger operators in 1D (and so on half line), Lieb–Thirring proved (initially for $p > \frac{1}{2}$, $p = \frac{1}{2}$ is Weidl and then Hundertmark–Lieb–Thomas)

$$\sum_{E_{j,\pm}} |E_j^\pm|^p \leq C_p \int_0^\infty |V(x)|^{p+\frac{1}{2}}$$

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Hundertmark–Simon (Killip–Simon for $p = \frac{3}{2}$)

$$\sum (|E_j^\pm| - 2)^p \leq \tilde{C}_p \sum_{n=0}^{\infty} |a_n - 1|^{p+\frac{1}{2}} + |b_n|^{p+\frac{1}{2}}$$

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The Gem

Hundertmark–Simon (Killip–Simon for $p = \frac{3}{2}$)

$$\sum (|E_j^\pm| - 2)^p \leq \tilde{C}_p \sum_{n=0}^{\infty} |a_n - 1|^{p+\frac{1}{2}} + |b_n|^{p+\frac{1}{2}}$$

Quasi-Szegő because power is $+\frac{1}{2}$, not $-\frac{1}{2}$ of Szegő condition.

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P_2 -Sum Rule

Define F on $\mathbb{R} \setminus [-2, 2]$ by ($|\beta| > 1$)

$$F(\beta + \beta^{-1}) = \frac{1}{4}[\beta^2 - \beta^{-2} - \log(\beta^4)];$$

$$F(E) = \frac{1}{2} \int_2^{|E|} (y^2 - 4)^{\frac{1}{2}} dy$$

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so $F(E) > 0$ and $F(E) = \frac{2}{3}(|E| - 2)^{\frac{3}{2}} + O((|E| - 2)^{\frac{5}{2}})$.

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$$Q(\mu) = \frac{1}{4\pi} \int_{-2}^2 \log\left(\frac{\sqrt{4-x^2}}{2\pi f(x)}\right) \sqrt{4-x^2} dx$$

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$$Q(\mu) = \frac{1}{4\pi} \int_{-2}^2 \log\left(\frac{\sqrt{4-x^2}}{2\pi f(x)}\right) \sqrt{4-x^2} dx$$

P_2 -Sum Rule:

$$Q(\mu) + \sum F(E_n^\pm) = \sum_{n=1}^{\infty} \left[\frac{1}{4} b_n^2 + \frac{1}{2} G(a_n) \right]$$

if $\sigma_{\text{ess}}(\mu) = [-2, 2]$

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P_2 -Sum Rule

$$\text{RHS} < \infty \Leftrightarrow \sum_{n=1}^{\infty} b_n^2 + (a_n - 1)^2 < \infty.$$

$$\text{LHS} < \infty \Leftrightarrow \text{Quasi-Szegő} + \sum_{n,\pm} (|E_n^\pm| - 2)^{\frac{3}{2}} < \infty.$$

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Thus P_2 -sum rule \Rightarrow KS Theorem.

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Case had formal sum rules depending on terms in a Taylor series.

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Case had formal sum rules depending on terms in a Taylor series. We called them C_n sum rules.

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Case had formal sum rules depending on terms in a Taylor series. We called them C_n sum rules. The P_2 sum rules is $C_0 + \frac{1}{2}C_2$.

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Case had formal sum rules depending on terms in a Taylor series. We called them C_n sum rules. The P_2 sum rule is $C_0 + \frac{1}{2}C_2$. It happens that due to mysterious cancellations, the terms in this combination are positive, which is why we used "P".

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Nazarov-Peherstorfer-Volberg-Yuditskii and then, Denisov Kupin for OPUC, did find ways of generating positive sum rules but they didn't seem to be computationally tractable.

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AMS

Real Analysis

Real Analysis
A Comprehensive Course in Analysis, Part 1

Barry Simon

$$xy \leq \frac{x^p}{p} + \frac{y^q}{q}$$

$$\hat{f}(\mathbf{k}) = (2\pi)^{-\nu/2} \int \exp(-i\mathbf{k} \cdot \mathbf{x}) f(\mathbf{x}) d^\nu x$$

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1

Simon

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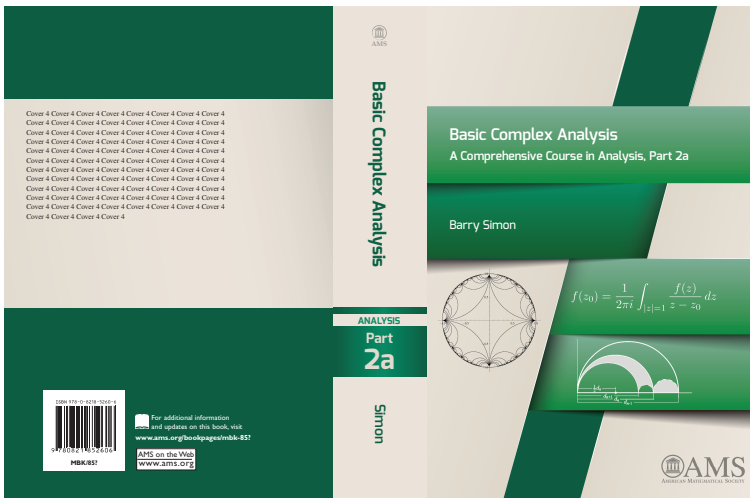
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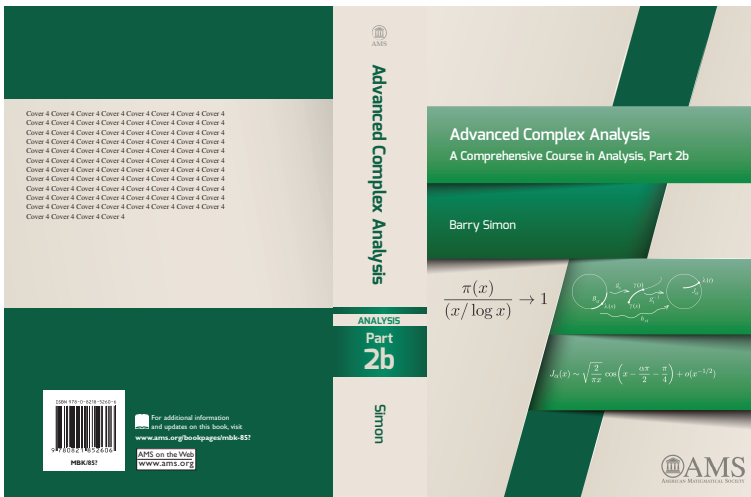
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