



Recent Developments in the Spectral Theory of Orthogonal Polynomials

Barry Simon

IBM Professor of Mathematics and Theoretical Physics
California Institute of Technology
Pasadena, CA, U.S.A.

Lecture 2: Sum Rules and Large Deviations

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- Lecture 2: Sum Rules and Large Deviations
- Lecture 3: Szegő-Widom asymptotics for Chebyshev Polynomials
- Lecture 4: Killip-Simon Theorems for Finite Gap Sets



References for Lecture 2

[GNR] F. Gamboa, J. Nagel, and A. Rouault
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[DS] J.D. Deuschel, and D. Stroock,
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[DZ] A. Dembo, and O. Zeitouni,
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The Punchline

First, I'm going to make propaganda for a year-old preprint of Gamboa, Nagel and Rouault (henceforth GNR) which has not been sufficiently appreciated.

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The Killip-Simon sum rule is just large deviations for GUE!!!!

Ironically, it appears the sum rules BSZ write using large deviations are just those of the Nazarov et al method that were previously regarded as untractable.



The LD Framework

Large deviations go back to Laplace. The modern theory was initiated by Cramer in the 1930's and made into a powerful machine by Donsker–Varadhan and then Varadhan alone (work for which he got the Abel prize).

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We consider a sequence of probability measures, $\{\mu_n\}_{n=1}^{\infty}$, on a space, X . Naively, one has a Large Deviation Principle (aka LDP) if the μ_n -probability that x is near x_0 is $O(e^{-nI(x_0)})$.

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We consider a sequence of probability measures, $\{\mu_n\}_{n=1}^{\infty}$, on a space, X . Naively, one has a Large Deviation Principle (aka LDP) if the μ_n -probability that x is near x_0 is $O(e^{-nI(x_0)})$. To be mathematically precise, one supposes that X is a Polish space (aka complete metric space), allows multiplicative factors other than n and so speaks of the *speed*, a_n , *rate function*, $I : X \rightarrow [0, \infty]$ and requires that:

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① I is lower semicontinuous

② For all closed sets $F \subset X$

$$\limsup_{n \rightarrow \infty} \frac{1}{a_n} \log \mu_n(F) \leq - \inf_{x \in F} I(x)$$

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$$I(x) = \sup_{\theta} [\theta x - \log(\mathbb{E}(e^{\theta X}))]$$

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LDP and Sum Rules

Gamboa, Nagel and Rouault had the following lovely idea.
Let X be the set of probability measures on $\partial\mathbb{D}$ or on \mathbb{R}
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LDP and Sum Rules

GNR had the further idea that the measures on the spectral measures should come from random matrix measures with a cyclic vector in the limit as the matrix dimension goes to infinity.

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LDP and Sum Rules

GNR had the further idea that the measures on the spectral measures should come from random matrix measures with a cyclic vector in the limit as the matrix dimension goes to infinity.

Of course, the issue becomes to effectively compute the rate function on both sides and alas, we haven't yet found a magic way to do these calculations in a general context.

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CUE: Measure Side

For CUE, we first consider Haar measure on $U(n)$, the $n \times n$ unitary matrices.

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CUE: Measure Side

For CUE, we first consider Haar measure on $U(n)$, the $n \times n$ unitary matrices. Any fixed vector is cyclic with probability one, so the corresponding spectral measures have the form $\sum_{j=1}^n w_j \delta_{\theta_j}$ where $\lambda_j \equiv e^{i\theta_j}$ are the eigenvalues.

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As is well-known, the λ 's and w 's are independent, the w 's are uniformly distributed on the simplex $\{\mathbf{w} \mid \sum_{j=1}^n w_j = 1\}$

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As is well-known, the λ 's and w 's are independent, the w 's are uniformly distributed on the simplex $\{\mathbf{w} \mid \sum_{j=1}^n w_j = 1\}$ and by the Weyl integration formula, the θ 's have distribution

$$\frac{1}{n!} \prod_{1 \leq j < k \leq n} |e^{i\theta_j} - e^{i\theta_k}|^2 \prod_{j=1}^n \frac{d\theta_j}{2\pi}$$

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CUE: Measure Side

The first step in the analysis of the measure side is to analyze what probabilists call the *empirical measure* and physicists *the density of states*, namely the random measure

$$\frac{1}{n} \sum_{j=1}^n \delta_{\theta_j}.$$

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The first step in the analysis of the measure side is to analyze what probabilists call the *empirical measure* and physicists *the density of states*, namely the random measure $\frac{1}{n} \sum_{j=1}^n \delta_{\theta_j}$. This also defines a family of measures on measures and, in 1997, Ben Arous and Guionnet made the important discovery that this (or rather an analog on the real line with a confining potential) has an LDP with speed n^2

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This is easy to understand.



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This is easy to understand. The Weyl distribution can be viewed as a discrete two dimensional Coulomb gas in the canonical ensemble



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This is easy to understand. The Weyl distribution can be viewed as a discrete two dimensional Coulomb gas in the canonical ensemble (2D because $|x - y|^{-2}$ is the exponential of $-2 \log |x - y|$). The $n \rightarrow \infty$ limit is a high density limit and due to repulsion, there is a strong tendency towards equal spacing.

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CUE: Measure Side

To get a significant difference from equal spacing, one has $O(n^2)$ smaller distances and so the speed is n^2 . The optimal spacing will still be locally equal and the discrete Coulomb energy will converge to the continuum.

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The fact that n^2 is much larger than n implies that for a measure to have finite rate at speed n , it has to have points close to uniformly distributed

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The weights are close to independent (except for the normalization they are)

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The weights are close to independent (except for the normalization they are) – a slick way to see this is to note if Y_j are positive exponentially distributed iidrv, then

$$w_j = Y_j / \sum_{j=1}^N Y_j.$$

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CUE: Measure Side

To get a significant difference from equal spacing, one has $O(n^2)$ smaller distances and so the speed is n^2 . The optimal spacing will still be locally equal and the discrete Coulomb energy will converge to the continuum.

The fact that n^2 is much larger than n implies that for a measure to have finite rate at speed n , it has to have points close to uniformly distributed and the large deviations comes from entirely from the lack of a uniform weight.

The weights are close to independent (except for the normalization they are) – a slick way to see this is to note if Y_j are positive exponentially distributed iidrv, then $w_j = Y_j / \sum_{j=1}^N Y_j$. This allows one (using the Chernoff-Cramer theorem on small blocks) to prove an LDP for the spectral measure with speed n and rate function the Szegő integral $-\int \log(w(\theta)) \frac{d\theta}{2\pi}$.

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CUE: Coefficient Side

In 2004, Killip and Nenciu wrote down the distribution of $\{\alpha_j\}_{j=0}^{n-1}$ induced by restricting Haar measure on a fixed vector as we are.

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$$\frac{n-j-1}{\pi} (1-|z|^2)^{n-j-2} d^2z$$

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$$\frac{n-j-1}{\pi} (1-|z|^2)^{n-j-2} d^2z$$

which says that α_j is distributed as the first complex component of a unit vector in \mathbb{C}^{n-j} .

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α_0 is U_{11} . Under Haar measure, each row is clearly uniformly distributed on the unit sphere in \mathbb{C}^n so we understand where the distribution of α_0 comes from.



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α_0 is U_{11} . Under Haar measure, each row is clearly uniformly distributed on the unit sphere in \mathbb{C}^n so we understand where the distribution of α_0 comes from. Intuitively, the rest of U which is a unitary map of δ_1^\perp to $U\delta_1^\perp$ is just Haar measure on this set of unitaries, so it is reasonable that the other α 's are independent and distributed according to CUE_{n-1} .



CUE: Coefficient Side

The Killip–Nenciu proof is more involved but there should be a proof along these lines.

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The Killip–Nenciu proof is more involved but there should be a proof along these lines.

$\prod \rho_j^2$ appears to the n th power so the rate function is $-\sum_{j=1}^{\infty} \log(1 - |\alpha_j|^2)$.

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My OPUC book has something like four other proofs of Szegő's Theorem, but until what I'll discuss next, there was really only one proof of the Killip-Simon theorem.

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GUE: Coefficient Side

The argument for GUE, normalized so the limiting density is the semicircle law on $[-2, 2]$, is similar.

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The argument for GUE, normalized so the limiting density is the semicircle law on $[-2, 2]$, is similar. Instead of results of Killip-Nenciu for the distribution of α 's, one has earlier results of Dumitriu and Edelman (2002) for the Jacobi parameters.

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GUE: Measure Side

One thing that this proof illuminates is why the Q term involves the quasi-Szegő $(4 - x^2)^{1/2}$ rather than the Szegő $(4 - x^2)^{-1/2}$.

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GUE: Measure Side

One thing that this proof illuminates is why the Q term involves the quasi-Szegő $(4 - x^2)^{1/2}$ rather than the Szegő $(4 - x^2)^{-1/2}$. The Szegő form is related to the equilibrium measure for $[-2, 2]$ which is the density of states for Jacobi matrices whose parameters go to the free ones.

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One needs to make some additional arguments going back to Ben Arous-Dembo-Guionnet (2001) to deal with eigenvalues outside the essential support. The positions of these eigenvalues matters and we get $F(E)$ terms.

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One needs to make some additional arguments going back to Ben Arous-Dembo-Guionnet (2001) to deal with eigenvalues outside the essential support. The positions of these eigenvalues matters and we get $F(E)$ terms.

What results is a new proof of the Killip-Simon sum rule.

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Higher Order Szegő Theorems

Killip-Simon was followed up by a variety of authors looking for other applications of sum rules. Among the authors who looked at this are Denisov, Golinskii, Kupin, Laptev, Lukic, Naboko, Nazarov, Novitskii, Peherstorfer, Safronov, Simon, Vainberg, Volberg, Yuditskii and Zlatoš.

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In writing my OPUC books, I realized that there was an analog of the combination of zeroth and second order that Killip-Simon used which led to the equality of

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In writing my OPUC books, I realized that there was an analog of the combination of zeroth and second order that Killip-Simon used which led to the equality of

$$\exp \left(\int_0^{2\pi} (1 - \cos \theta) \log w(\theta) \frac{d\theta}{2\pi} \right)$$

and

$$\exp \left(\frac{1}{2} - \frac{1}{2} |\alpha_0 + 1|^2 \right) \prod_{n=0}^{\infty} e^{-\frac{1}{2} |\alpha_{n+1} - \alpha_n|^2} \prod_{n=0}^{\infty} (1 - |\alpha_n|^2) e^{|\alpha_n|^2}$$

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Notice the $|\alpha_0 + 1|^2$ term.

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Notice the $|\alpha_0 + 1|^2$ term. It is irrelevant for deducing gems and is a finite boundary term.

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Notice the $|\alpha_0 + 1|^2$ term. It is irrelevant for deducing gems and is a finite boundary term. As we go further, these terms become more involved and we will throw them in an “always finite” basket and ignore them.

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Because the last exponential cancels the first term in the expansion of $\log(1 - |\alpha_n|^2)$, we get the gem

$$\exp \left(\int_0^{2\pi} (1 - \cos \theta) \log w(\theta) \frac{d\theta}{2\pi} \right) > -\infty \iff \sum_{n=0}^{\infty} (|\alpha_{n+1} - \alpha_n|^2 + |\alpha_n|^4) < \infty$$

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Basically, w is allowed to have a higher order zero at $\theta = 0$ but one loses the ℓ^2 property of the α 's which are allowed to decay more slowly.

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Simon's Conjecture

Motivated by this example, I made the following conjecture:

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Simon's Conjecture

Motivated by this example, I made the following conjecture:
Let $\theta_1, \dots, \theta_k \in [0, 2\pi)$ be distinct and m_1, \dots, m_k strictly positive integers, let $q = 1 + \max_{j=1, \dots, k} m_j$ and let S be the operator $(S\alpha)_k = \alpha_{k+1}$. Then

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$$\int \prod_{j=1}^k [1 - \cos(\theta - \theta_j)]^{m_j} \log w(\theta) \frac{d\theta}{2\pi} > -\infty \iff$$

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$$\int \prod_{j=1}^k [1 - \cos(\theta - \theta_j)]^{m_j} \log w(\theta) \frac{d\theta}{2\pi} > -\infty \iff$$

$$\prod_{j=1}^k (S - e^{-i\theta_j})^{m_j} \alpha \in \ell^2 \quad \text{and} \quad \alpha \in \ell^{2q}$$

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$$\prod_{j=1}^k (S - e^{-i\theta_j})^{m_j} \alpha \in \ell^2 \quad \text{and} \quad \alpha \in \ell^{2q}$$

Above, we discussed the case $k = 1, m_1 = 1$. Simon-Zlotos proved this for $k = 1, m_1 = 2$ and $k = 2, m_1 = m_2 = 1$.

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Lukic's Conjecture

By using a flawed gem, Lukic was able to find a counterexample to the next case, viz: $k = 2, \theta_1 = 0, \theta_2 = \pi, m_2 = 2, m_1 = 1!!!$

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But he came up with a modified conjecture:

$$\int_0^{2\pi} \left[\prod_{j=1}^n (1 - \cos(\theta - \theta_j))^{m_j} \right] \log(w(\theta)) d\theta > -\infty \iff$$

$$\alpha = \beta^{(1)} + \dots + \beta^{(n)} \quad (S - e^{-i\theta_j})^{m_j} \beta^{(j)} \in \ell^2 \quad \beta^{(j)} \in \ell^{2m_j+2}$$

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But he came up with a modified conjecture:

$$\int_0^{2\pi} \left[\prod_{j=1}^n (1 - \cos(\theta - \theta_j))^{m_j} \right] \log(w(\theta)) d\theta > -\infty \iff$$

$$\alpha = \beta^{(1)} + \dots + \beta^{(n)} \quad (S - e^{-i\theta_j})^{m_j} \beta^{(j)} \in \ell^2 \quad \beta^{(j)} \in \ell^{2m_j+2}$$

In some sense the α 's are non-local moments, aka Fourier coefficients, so the measure is sort of the non-linear Fourier sum of α 's,

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Lukic's Conjecture

By using a flawed gem, Lukic was able to find a counterexample to the next case, viz: $k = 2, \theta_1 = 0, \theta_2 = \pi, m_2 = 2, m_1 = 1!!!$

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In some sense the α 's are non-local moments, aka Fourier coefficients, so the measure is sort of the non-linear Fourier sum of α 's, so the β 's are the localized pieces.

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Lukic's Conjecture

While the decomposition is conceptually useful, it is not clear how to realize it as part of a sum rule,

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The project that BSZ has in process is to use the insights of GNR to approach Lukic's conjecture. We are focusing, for now, on the cases where $k \leq 3, \theta_1 = 0, \theta_2 = \pi$.

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The project that BSZ has in process is to use the insights of GNR to approach Lukic's conjecture. We are focusing, for now, on the cases where $k \leq 3, \theta_1 = 0, \theta_2 = \pi$. We have recovered the known cases where (m_1, m_2) is $(1, 0)$ or $(1, 1)$ or $(2, 0)$.

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The project that BSZ has in process is to use the insights of GNR to approach Lukic's conjecture. We are focusing, for now, on the cases where $k \leq 3, \theta_1 = 0, \theta_2 = \pi$. We have recovered the known cases where (m_1, m_2) is $(1, 0)$ or $(1, 1)$ or $(2, 0)$. Our next goal will be to get $(2, 1)$ where Lukic's conjecture is different from my (incorrect) conjecture. To the extent that time allows, I'll say something about the details of our calculations.

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Measure Side

One thing that made the calculations for CUE and GUE easy was independence of almost everything in sight (eigenvalues, α 's, and almost the weights).

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Measure Side

One thing that made the calculations for CUE and GUE easy was independence of almost everything in sight (eigenvalues, α 's, and almost the weights). The occurrence of difference operators says we'll lose independence but, by looking at multiplicative perturbations of CUE, that will be tractable.

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One thing that made the calculations for CUE and GUE easy was independence of almost everything in sight (eigenvalues, α 's, and almost the weights). The occurrence of difference operators says we'll lose independence but, by looking at multiplicative perturbations of CUE, that will be tractable. So we'll look at measures on random unitary matrices of the form

$$d\nu_n = Z_n^{-1} \exp\left(-n \sum_{j=1}^n V(\lambda_j)\right) d\nu_n^{(0)}$$

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$$d\nu_n = Z_n^{-1} \exp\left(-n \sum_{j=1}^n V(\lambda_j)\right) d\nu_n^{(0)}$$

where $d\nu_n^{(0)}$ is CUE, λ_j are the eigenvalues and Z_n is a normalizing factor.

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Measure Side

We are interested in getting a limiting empirical measure of the form $d\eta(\theta)$.

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Measure Side

We are interested in getting a limiting empirical measure of the form $d\eta(\theta)$. Prior work says that we get this if V is picked to be (confusing $\lambda = e^{i\theta}$ and θ):

$$V(\theta) = 2 \int \log |e^{i\theta} - e^{i\psi}| d\eta(\psi)$$

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In that case, the rate function for the spectral measure is just a relative entropy,

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In that case, the rate function for the spectral measure is just a relative entropy, so if $d\eta$ is absolutely continuous, up to an additive constant, the rate function is $-\int \log w(\theta) d\eta(\theta)$.

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In that case, the rate function for the spectral measure is just a relative entropy, so if $d\eta$ is absolutely continuous, up to an additive constant, the rate function is $-\int \log w(\theta) d\eta(\theta)$. Thus, to explore cases of Lukic's conjecture, we need to compute

$$V(\theta) = 2 \frac{\int \log |e^{i\theta} - e^{i\psi}| \left[\prod_{j=1}^n (1 - \cos(\psi - \theta_j))^{m_j} \right] d\psi}{\int \left[\prod_{j=1}^n (1 - \cos(\psi - \theta_j))^{m_j} \right] d\psi}$$

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Measure Side

It is useful to know that

$$\begin{aligned}\int_0^{2\pi} \log[1 - \cos(\psi)] \cos(n\psi) \frac{d\psi}{2\pi} &= \int_0^{2\pi} \log[1 - \cos(\psi)] e^{in\psi} \frac{d\psi}{2\pi} \\ &= \frac{1}{n}\end{aligned}$$

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and that

$$|1 - e^{i\theta}|^2 = 2(1 - \cos(\theta))$$

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and that

$$|1 - e^{i\theta}|^2 = 2(1 - \cos(\theta))$$

This is used inside a log so the 2 gives an irrelevant additive constant but the square gives a multiplicative factor of 2 which is important.

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and that

$$|1 - e^{i\theta}|^2 = 2(1 - \cos(\theta))$$

This is used inside a log so the 2 gives an irrelevant additive constant but the square gives a multiplicative factor of 2 which is important. These formulae show that for the (1,0) case, $V(\theta) = \cos(\theta)$.

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Coefficient Side

$$\cos(\theta) = \operatorname{Re}(e^{i\theta}), \text{ so } \sum V(\lambda_j) = \operatorname{Re}(\operatorname{Tr}(U)).$$

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Coefficient Side

$\cos(\theta) = \operatorname{Re}(e^{i\theta})$, so $\sum V(\lambda_j) = \operatorname{Re}(\operatorname{Tr}(U))$. The trace can be computed in any basis and it is convenient to pick the basis in which U is a so-called CMV matrix which is five diagonal with $(\bar{\alpha}_0, -\bar{\alpha}_1\alpha_0, \dots, -\bar{\alpha}_{n+1}\alpha_n, \dots)$ along the diagonal so

$$\operatorname{Re}(\operatorname{Tr}(U)) = \frac{1}{2} \left[\bar{\alpha}_0 + \alpha_0 - \sum_{n=0}^{\infty} (\bar{\alpha}_{n+1}\alpha_n + \alpha_{n+1}\bar{\alpha}_n) \right]$$

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Coefficient Side

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The CUE rate function has this added to it, so completing the square, one recovers the sum rule that I used to do the $(1, 0)$ case.

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Coefficient Side

$\cos(\theta) = \text{Re}(e^{i\theta})$, so $\sum V(\lambda_j) = \text{Re}(\text{Tr}(U))$. The trace can be computed in any basis and it is convenient to pick the basis in which U is a so-called CMV matrix which is five diagonal with $(\bar{\alpha}_0, -\bar{\alpha}_1\alpha_0, \dots, -\bar{\alpha}_{n+1}\alpha_n, \dots)$ along the diagonal so

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The CUE rate function has this added to it, so completing the square, one recovers the sum rule that I used to do the (1,0) case. What is striking is that I had to be clever to get the exact sum rule while here is is automated.

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Measure Side

The potential is, up to an additive constant (which doesn't matter for gems), $\frac{1}{2} \cos(2\theta)$.

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Measure Side

The potential is, up to an additive constant (which doesn't matter for gems), $\frac{1}{2} \cos(2\theta)$. In this calculation, one uses that the normalized measure (which for the $(1, 0)$ case was $d\eta = (2\pi)^{-1}(1 - \cos(\theta))d\theta$) is now $d\eta = \pi^{-1}(1 - \cos^2(\theta))d\theta$.

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Coefficient Side

$\cos(2\theta)$ looks like an innocuous 2 but on the coefficient side, it corresponds to $\text{ReTr}(U^2)$ which is more complicated than $\text{Tr}(U)$ since U is 5 diagonal.

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Coefficient Side

$\cos(2\theta)$ looks like an innocuous 2 but on the coefficient side, it corresponds to $\text{ReTr}(U^2)$ which is more complicated than $\text{Tr}(U)$ since U is 5 diagonal. It turns out that there are only three terms and two are equal, so up to finite boundary terms,

$$\text{Tr}(U^2) = \sum_{j=0}^{\infty} (\bar{\alpha}_j \alpha_{j+1})^2 - 2 \sum_{j=0}^{\infty} \rho_j^2 \bar{\alpha}_{j+1} \alpha_{j-1}$$

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One expands ρ and also the sum of logs and combines all second order terms to get that the rate function is (up to finite terms)

$$\frac{1}{2} \left[\sum_{j=0}^{\infty} |\alpha_{j+1} - \alpha_{j-1}|^2 + E_1 + E_2 \right] + \sum_{k=2}^{\infty} k^{-1} \sum_{j=0}^{\infty} |\alpha_j|^{2k}$$

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Coefficient Side

where

$$E_1 = \sum_{j=0}^{\infty} (|\alpha_j|^2 |\alpha_{j+1}|^2 + \operatorname{Re}[(\bar{\alpha}_j \alpha_{j+1})^2])$$

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Coefficient Side

where

$$E_1 = \sum_{j=0}^{\infty} (|\alpha_j|^2 |\alpha_{j+1}|^2 + \operatorname{Re}[(\bar{\alpha}_j \alpha_{j+1})^2])$$
$$E_2 = \sum_{j=1}^{\infty} |\alpha_j|^2 \operatorname{Re}[\bar{\alpha}_j (\alpha_{j-1} - \alpha_{j+1})]$$

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Coefficient Side

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$$E_2 = \sum_{j=1}^{\infty} |\alpha_j|^2 \operatorname{Re}[\bar{\alpha}_j (\alpha_{j-1} - \alpha_{j+1})]$$

The goal is to prove that:

$$\sum_{j=1}^{\infty} |\alpha_{j+1} - \alpha_{j-1}|^2 + |\alpha_j|^4 < \infty \iff \text{rate function finite}$$

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Coefficient Side

where

$$E_1 = \sum_{j=0}^{\infty} (|\alpha_j|^2 |\alpha_{j+1}|^2 + \operatorname{Re}[(\bar{\alpha}_j \alpha_{j+1})^2])$$

$$E_2 = \sum_{j=1}^{\infty} |\alpha_j|^2 \operatorname{Re}[\bar{\alpha}_j (\alpha_{j-1} - \alpha_{j+1})]$$

The goal is to prove that:

$$\sum_{j=1}^{\infty} |\alpha_{j+1} - \alpha_{j-1}|^2 + |\alpha_j|^4 < \infty \iff \text{rate function finite}$$

Since E_1 and E_2 are fourth order, LHS and Hölder's inequality proves that the LHS \Rightarrow RHS.

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Coefficient Side

On the other hand, suppose that the rate function is finite.

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Coefficient Side

On the other hand, suppose that the rate function is finite. By using the measure side, we see that the a.c. part of the spectrum is the entire unit circle, so by a Theorem of Rakhmanov, $\alpha_j \rightarrow 0$.

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Coefficient Side

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$$|E_2| \leq \left[\sum_{j=1}^{\infty} |\alpha_{j-1} - \alpha_{j+1}|^2 \right]^{1/2} \left[\sum_{j=1}^{\infty} |\alpha_j|^6 \right]^{1/2}$$

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$$|E_2| \leq \left[\sum_{j=1}^{\infty} |\alpha_{j-1} - \alpha_{j+1}|^2 \right]^{1/2} \left[\sum_{j=1}^{\infty} |\alpha_j|^6 \right]^{1/2}$$

Since $\alpha_j \rightarrow 0$, this can be controlled by a small amount of $\sum_{j=1}^{\infty} |\alpha_{j+1} - \alpha_{j-1}|^2 + |\alpha_j|^4$ so we conclude that this latter sum is finite showing that $\text{RHS} \Rightarrow \text{LHS}$.

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Coefficient Side

This bare hands approach gets harder and harder as orders increase but we think we have $(2, 0)$ (where $V(\theta) = \frac{4}{3} \cos(\theta) - \frac{1}{6} \cos(2\theta) + \text{constant}$) and hope to do $(2, 1)$.

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Coefficient Side

This bare hands approach gets harder and harder as orders increase but we think we have $(2, 0)$ (where $V(\theta) = \frac{4}{3} \cos(\theta) - \frac{1}{6} \cos(2\theta) + \text{constant}$) and hope to do $(2, 1)$. It also is plausible there is some clever way of avoiding too many explicit calculations.

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$$xy \leq \frac{x^p}{p} + \frac{y^q}{q}$$



$$\hat{f}(\mathbf{k}) = (2\pi)^{-\nu/2} \int \exp(-i\mathbf{k} \cdot \mathbf{x}) f(\mathbf{x}) d^\nu x$$

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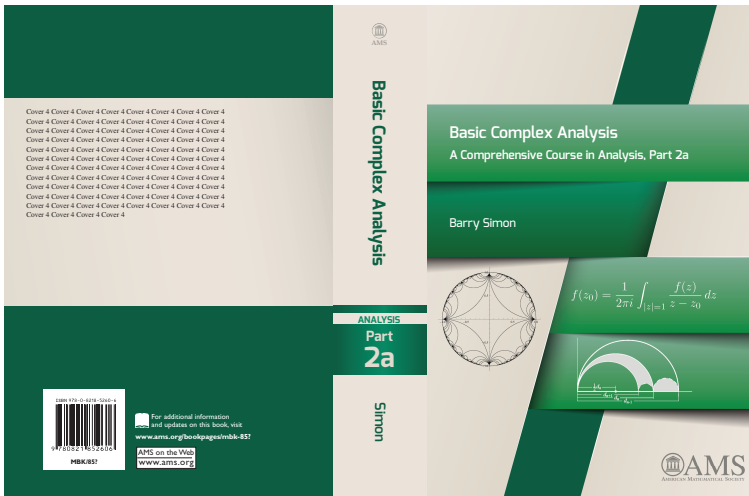
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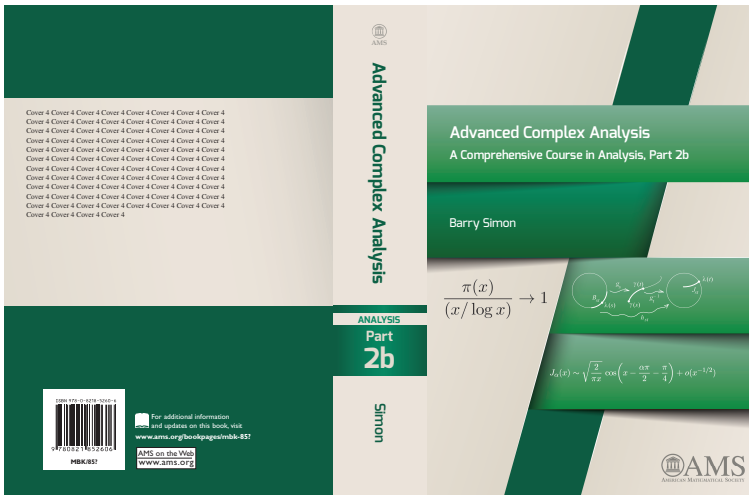
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